

# LASER-INDUCED NANOPARTICLE FORMATION IN LIQUIDS: INVOLVED MECHANISMS AND ROLE OF THE EXPERIMENTAL PARAMETERS

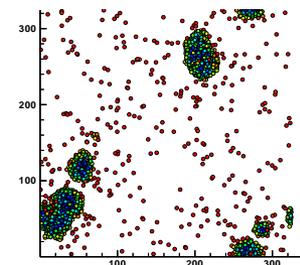
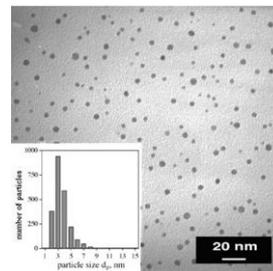
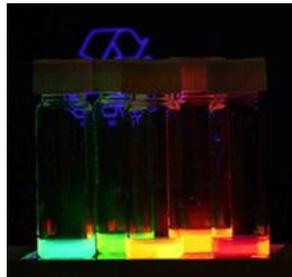
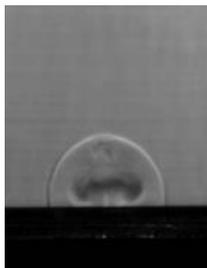
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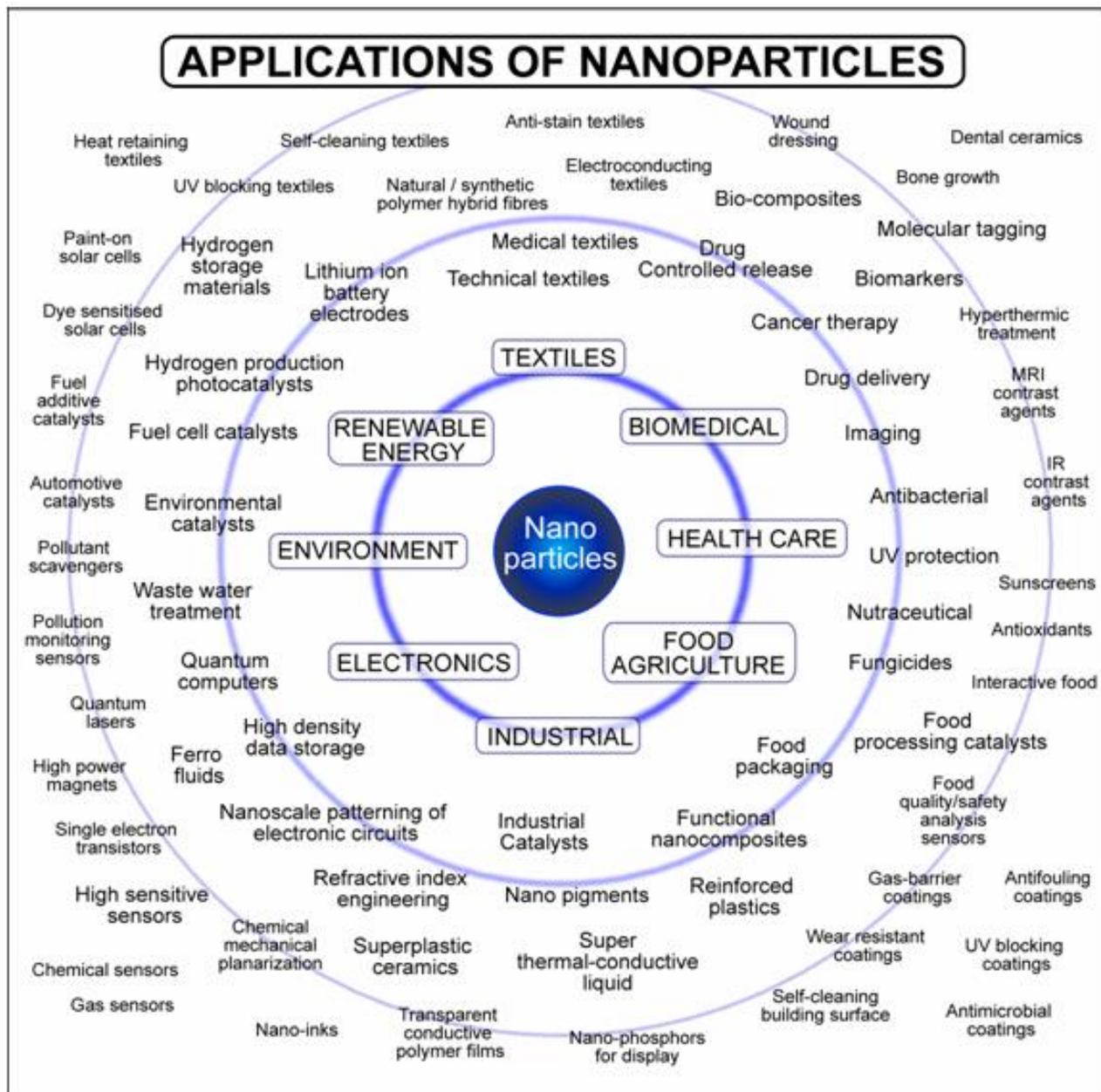


# OUTLINE

- 1. Importance and motivation**
- 2. Laser ablation in vacuum and in a background gas**
  - Shock waves (SW) and NP formation
- 3. Laser ablation in liquids**
  - “Primary” particles: direct ejection, nucleation in bubble or in solution ?
  - Secondary” particles: evaporation/growth, coalescence, agglomeration
- 4. Laser-induced particle size reduction and fragmentation**
  - Role of laser absorption by particles, fragmentation, stabilization
- 5. Summary**

# ***1. Importance and motivation***

# Motivation



# Applications of Metallic Nanoparticles (NPs)

## *Plasmonic properties*

-Sensors

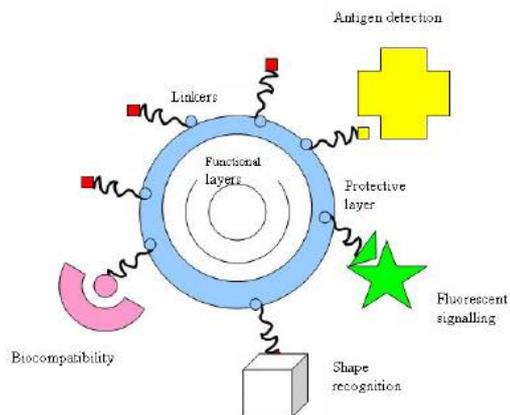
-New generation of solar cells

-**Cancer treatment** (photodynamic therapy, photo-thermal therapy, radiation therapy, drug delivery, etc...)

-**Diagnostic imaging** MRI (contrast agents), fluorescence Imaging, optical coherence tomography, opto-acoustic tomography, radionuclide Imaging, ultrasound imaging

-**Antiseptics and tissue engineering**

**Green Chemistry ?** *Laser-produced colloidal NPs are chemically clean and biocompatible*

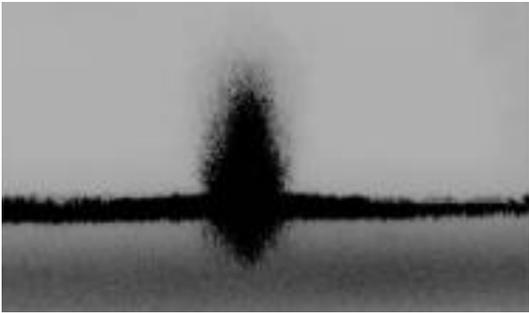
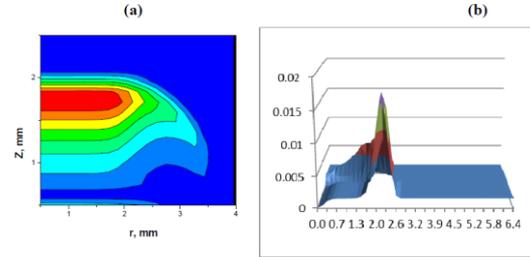
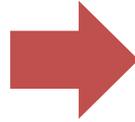
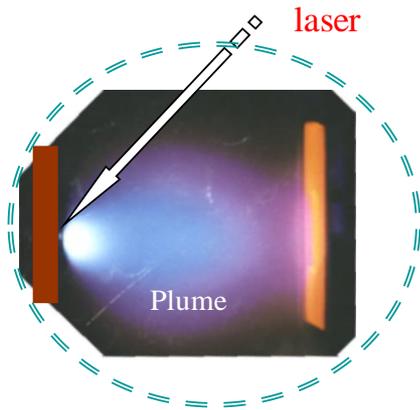


• *Photon-based nanoscience and nanobiotechnology, Nato Science Series Vol. 239*

• *Journal of NanoBiotechnology*

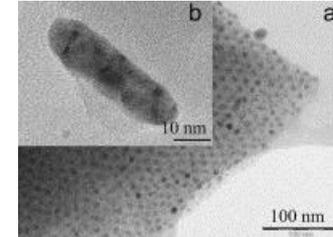
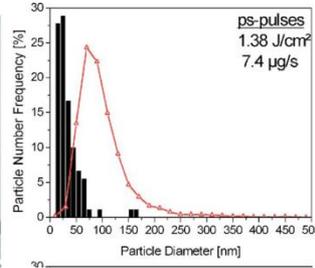
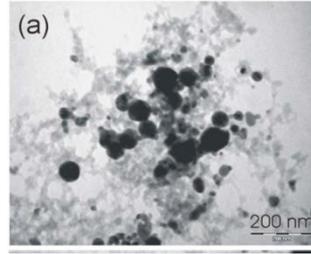
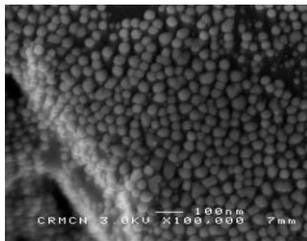
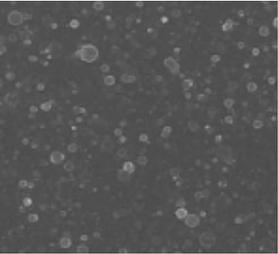
*Salata, 2004*

# Laser-assisted NP formation



- Laser beam is focused on a target
- Target material is heated and ejected
- Nanoparticles are formed

Heiroth et al. JAP 2009



Deposited nanoparticles,  
Courtesy of  
Garrelie et al.

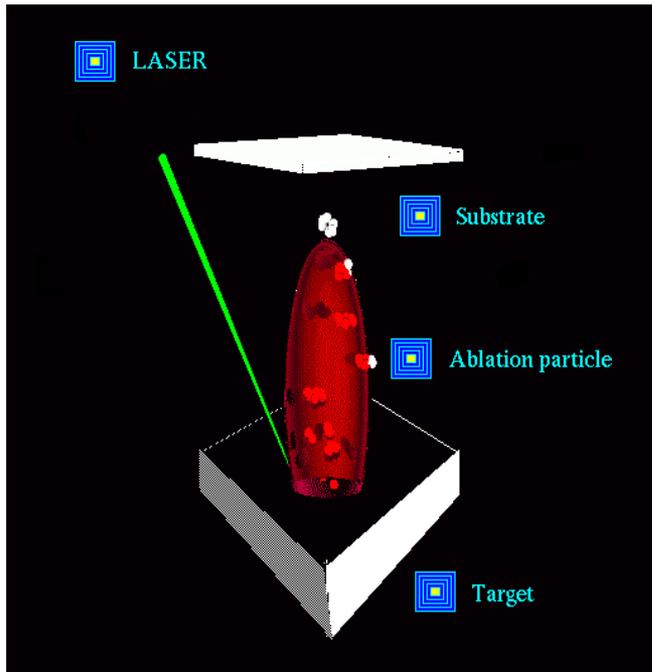
Nanoparticles produced  
on the target  
Kabashin et al.

Colloidal nanoparticles  
Barcikowski et al. APL 2007

Sol-gel nanoparticle arrays  
Bois et al.  
J. Sol. State Chem. 2009

## ***2. Laser ablation in vacuum and in a background gas***

# Processes involved



- Radiation absorption by target and by plasma
- Creation of regions with high T et P
- Propagation of pressure and thermal waves
- Phase transitions and material decomposition
- Ionization
- Ejection of electrons , ions and particles
- Plasma plume formation
- Collisions and chemical reactions
- Formation of clusters and nanoparticles (NPs)
- Emission
- More ...

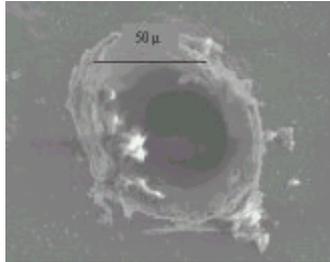
⇒ *Many physico-chemical processes*

⇒ *Depend on target material and laser source*

⇒ *Different time scales*

# Role of Pulse Duration

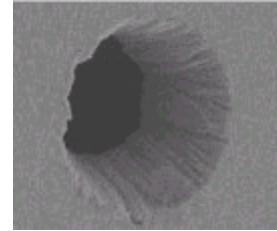
## NANOSECOND



- heat diffusion
- melting
- boiling
- evaporation

$$J \sim \frac{p_b}{(2\pi k T m)^{1/2}} \exp \left\{ \frac{\Delta H_v(T_b) m}{k} \left( \frac{1}{T_b} - \frac{1}{T} \right) \right\}$$

## FEMTOSECOND



- $\tau$  is shorter than  $\tau_{ei}$
- thermal conductivity during the pulse is negligible
- strong pressure gradients
- less energy is transformed into heat and more into motion



**Ablation mechanisms ?**

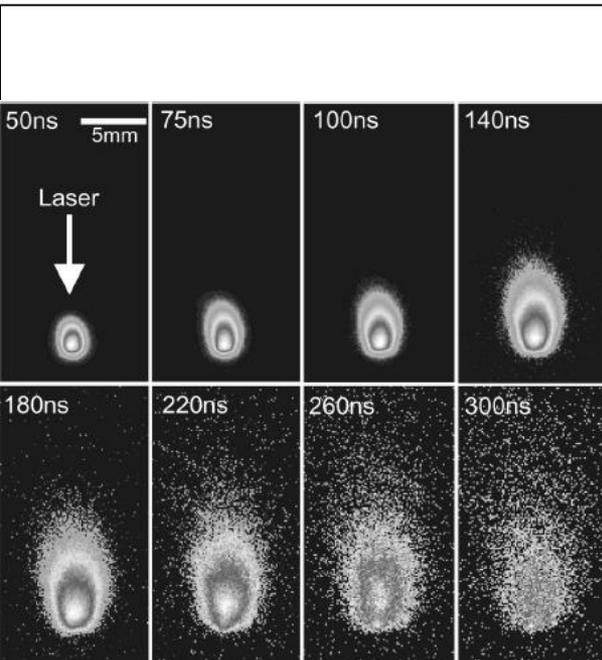
# Nanosecond Laser Pulses

In vacuum

In a gas

Low pressure

Atmospheric pressure



J. Appl. Phys., Vol. 93, No. 5, 1 March 2003

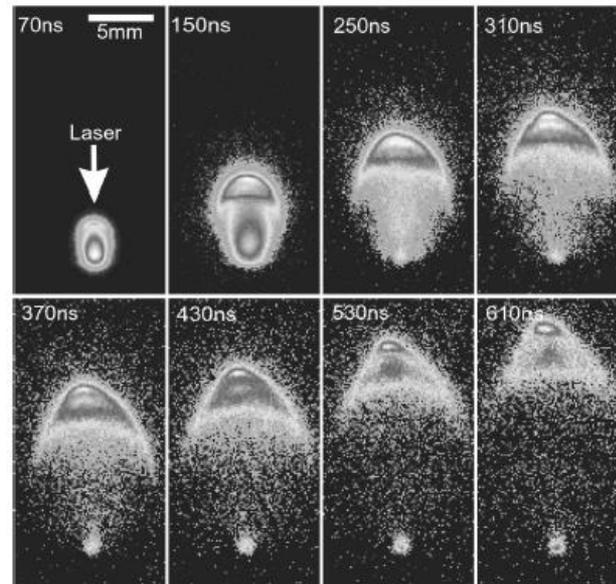
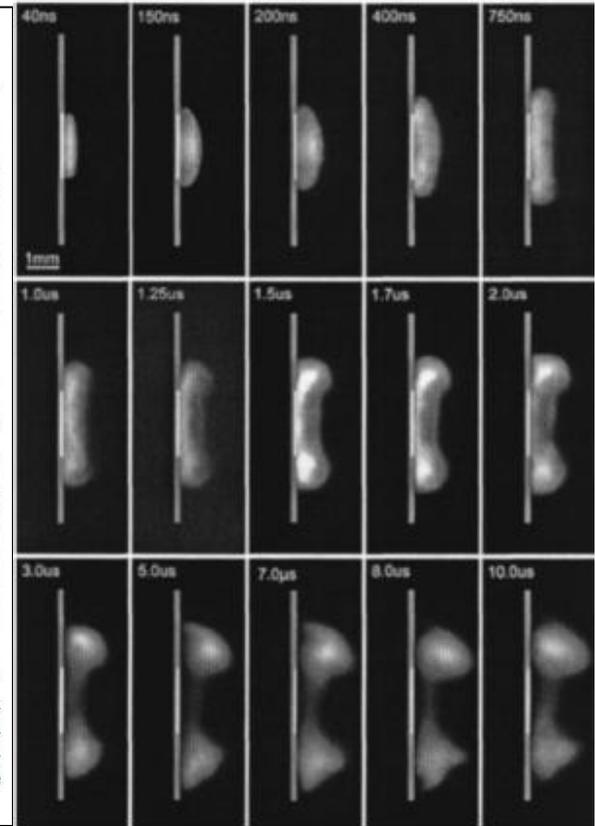


FIG. 4. ICCD photographs of visible emission from laser-produced aluminum plasma at 150 mTorr background air pressure. The experimental parameters are the same as described in Fig. 2. All of the images are normalized to their maximum intensity. Plume splitting and sharpening are observed in this pressure regime.



Pereira *et al.* XeCl, 30 ns, 10 J/cm<sup>2</sup>, steel

J. Appl. Phys. 98, 064902 (2005)

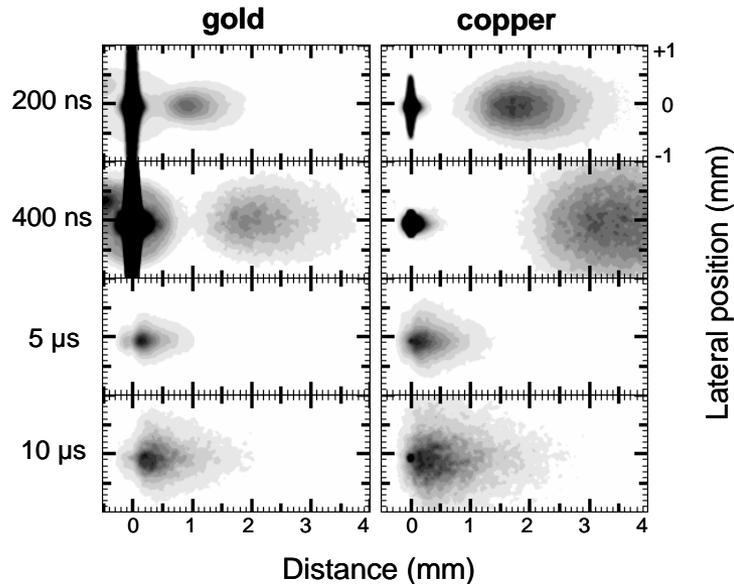
Harilal *et al.*, JAP 2003

Nd/YAG laser (8 ns, 700 mJ, Al target)

Stoichiometry is better preserved than in SD!

# Femtosecond Laser Pulses

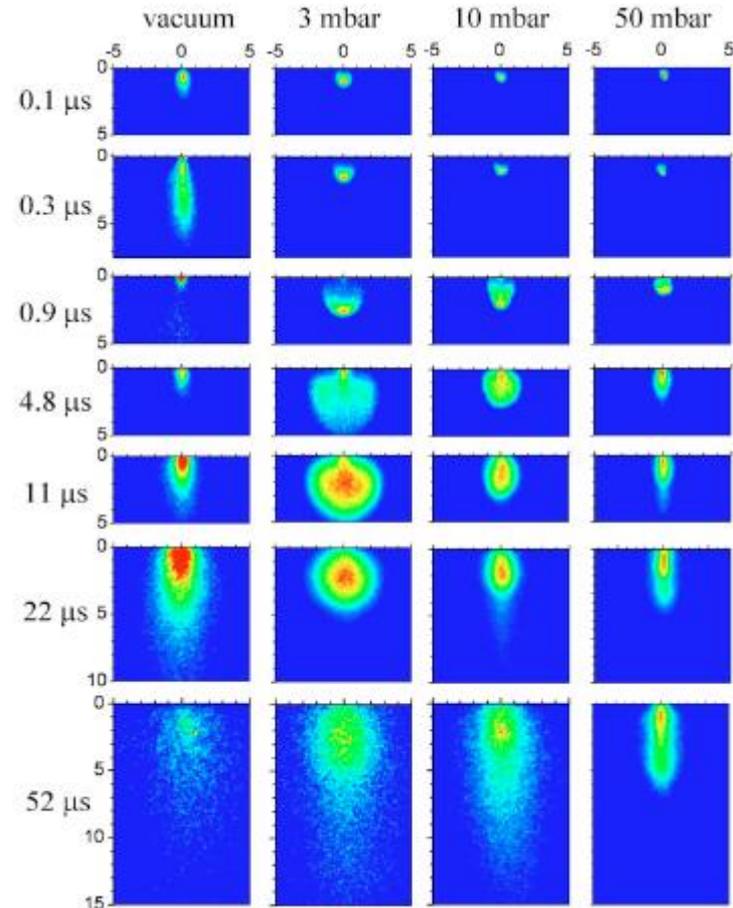
In vacuum



Plume images recorded during ablation of gold and copper for several times and a laser fluence of  $4 \text{ J cm}^{-2}$

*Hermann et al., Laser Phys. 2008*

In a gas



Plume images recorded during ablation of Fe in a background gas

*Amoruso et al, APL 2008*

## Two components were also observed

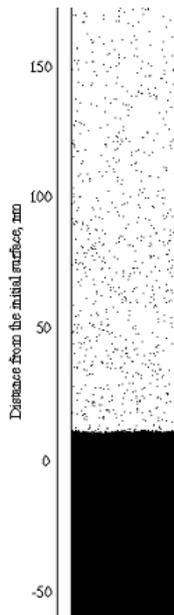
- by Garrelie et al. for fs ablation of carbon in vacuum (2001)
- By Grojo et al for fs ablation of Ti, Zr, Hf in vacuum (2002-2003)

# NP's formation in vacuum

*Molecular dynamics simulations*

Longer laser pulses (>ns)  
Small laser intensity

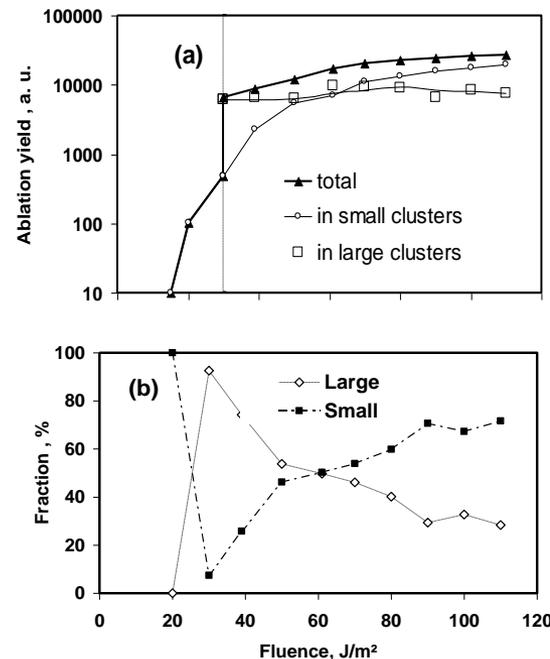
Laser



Target

+ nucleation/growth in the plume via condensation, coalescence and/or aggregation

Short laser pulses (ps, fs)  
High laser intensity



⇒ direct ejection of nanoclusters

⇒ the higher is F the smaller are the particles

# Two temperature hydrodynamics

$$\frac{\partial V}{\partial t} - \frac{\partial u}{\partial m} = 0, \quad (1)$$

$$\frac{\partial u}{\partial t} + \frac{\partial(P_i + P_e)}{\partial m} = 0, \quad (2)$$

$$\frac{\partial e_e}{\partial t} + P_e \frac{\partial u}{\partial m} = -\gamma_{ei}(T_e - T_i)V + Q_L V + \frac{\partial}{\partial m} \left( \kappa \frac{\partial T_e}{\partial z} \right), \quad (3)$$

$$\frac{\partial e_i}{\partial t} + P_i \frac{\partial u}{\partial m} = \gamma_{ei}(T_e - T_i)V. \quad (4)$$

- 2T Lagrangian hydrodynamics
- 2T semi-empirical equations of state (for metal, gold here)
- Electron-ion coupling
- Thermal conductivity
- Wide range permittivity

$$\nu_{\text{eff}} = \min(\nu_{\text{met}}, \nu_{\text{max}}, \nu_{\text{pl}}) \quad \gamma_{ei} = \frac{3k_B m_e}{m_i} n_e \nu_{\text{eff}}$$

$$\kappa_{\text{met}} = \frac{\pi^2 k_B^2 n_e}{3m_e \nu_{\text{eff},t}} T_e \quad T_e \ll T_F$$

$$\kappa = \kappa_{\text{pl}} + (\kappa_{\text{met}} - \kappa_{\text{pl}}) e^{-A_4^i T_e / T_F}$$

$$\kappa_{\text{pl}} = \frac{16 \sqrt{2} k_B (k_B T_e)^{5/2}}{\pi^{3/2} Z e^4 \sqrt{m_e} \Lambda} \quad T_e \gg T_F$$

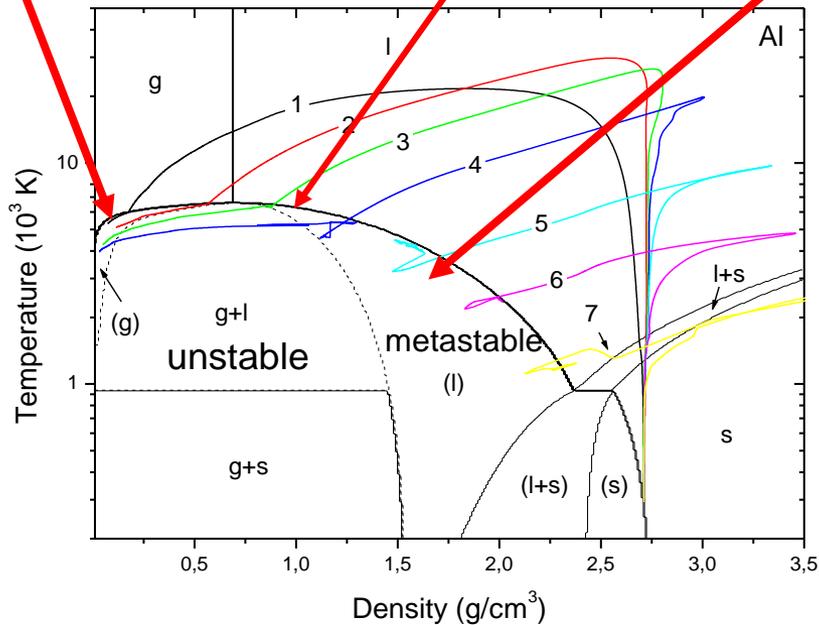
$$\varepsilon = \varepsilon_{\text{pl}} + (\varepsilon_{\text{met}} - \varepsilon_{\text{pl}}) e^{-A_4^p T_e / T_F}$$

# Thermodynamic analysis of laser ablation in vacuum

3-Nucleation and condensation

1-Phase explosion

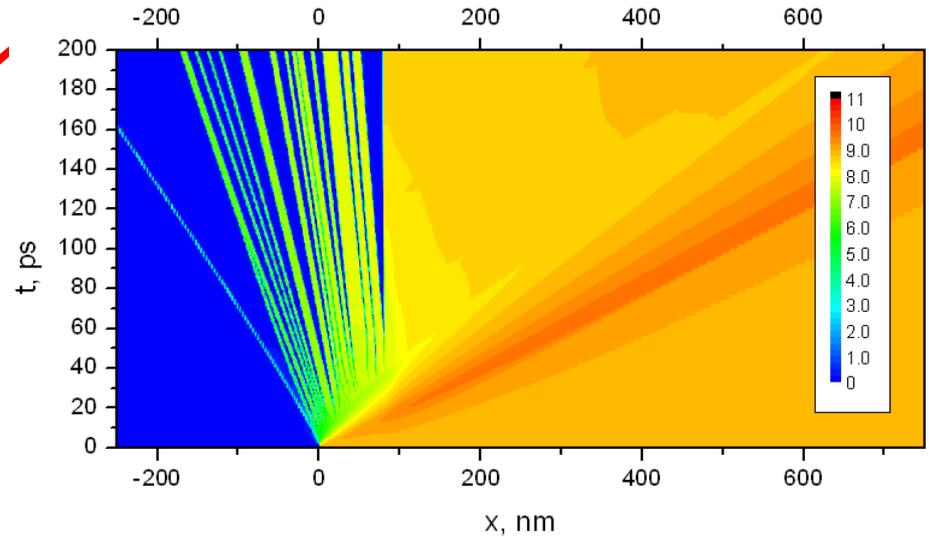
$$0.9T_c < T < T_c$$



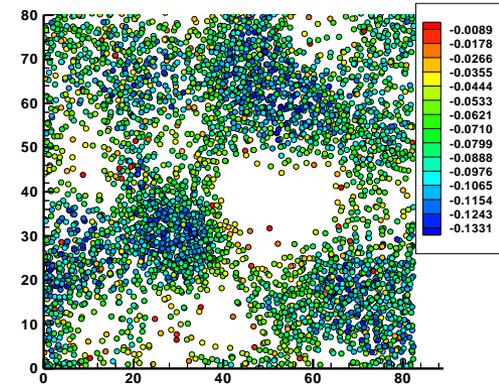
$$\tau_L = 100 \text{ fs}, \lambda = 800 \text{ nm}, F = 5.0 \text{ J/cm}^2$$

- In vacuum, fragmentation of metastable liquid leads to the ejection of clusters and chunks
- What is the mechanism in liquid ?

2- Fragmentation

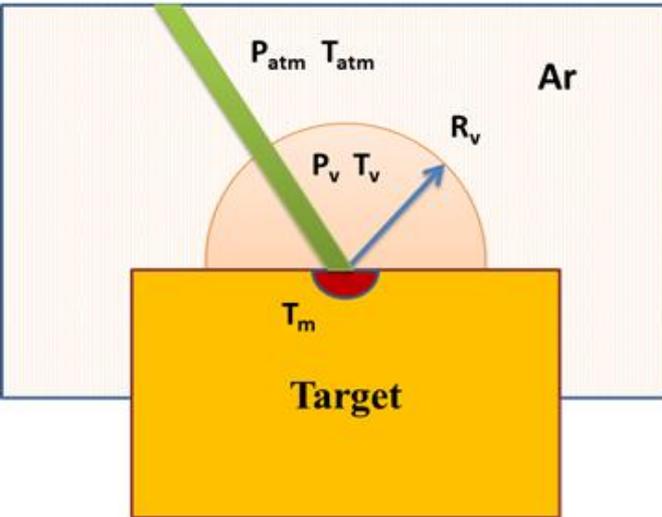


*Phys. Rev. B*, 75, 235414 (2007)



*Chem. Phys. Lett.* (2008)

# In a gas: Shock wave



**Shock wave (SW)**

**Estimation : 3D blast wave**

$$R = \xi \left( \frac{2E_0}{\rho} \right)^{1/5} t^{2/5}$$

$$\frac{T}{T_0} = \left( \frac{V_0}{V} \right)^{\gamma-1} \Rightarrow T = T_0 \left( \frac{R_0}{R} \right)^{3(\gamma-1)}$$

$$T_{eq} = \left\{ \frac{1}{T_b} - \frac{k}{Q} \ln \left( \frac{P}{P_0} \right) \right\}^{-1}$$

$$\theta = \frac{T_{eq} - T}{T_{eq}}, \text{ see Zel'dovich \& Raizer}$$

**Supersaturation degree vs time:**

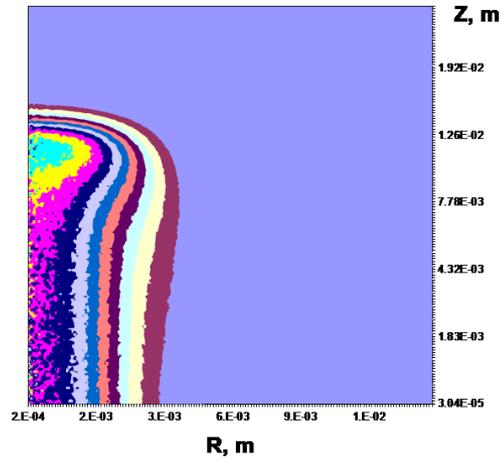
$$r^* = \frac{2\sigma v}{\theta k Q}$$

$$\theta = 1 - \frac{T}{T_{eq}} = 1 - T \left\{ \frac{1}{T_b} - \frac{k}{Q} \ln \left( \frac{P}{P_{atm}} \right) \right\} =$$

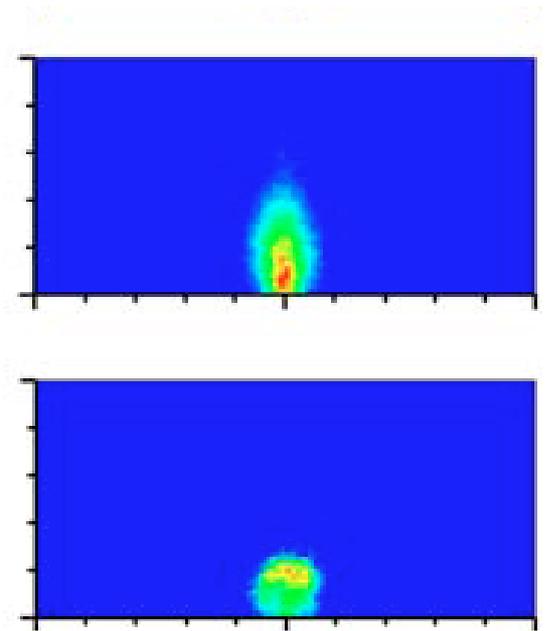
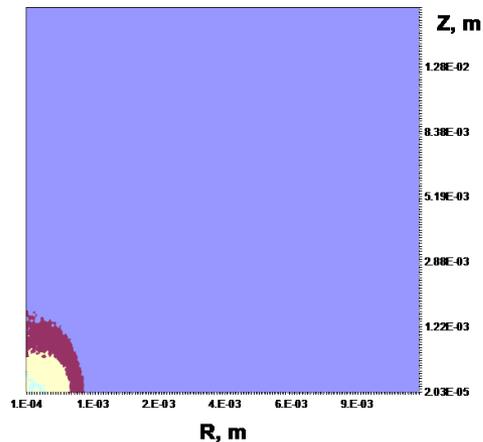
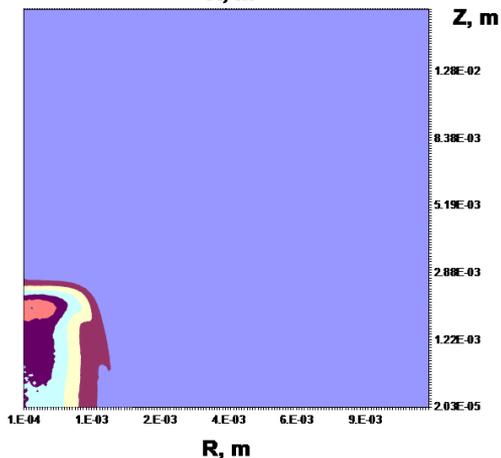
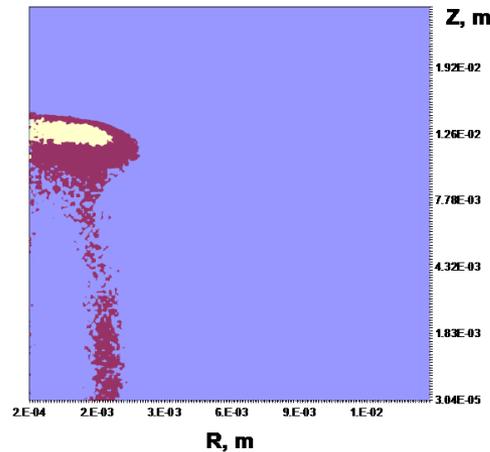
$$1 - \frac{T_0}{T_b} \left( \frac{R_0}{R} \right)^{3(\gamma-1)} \left\{ 1 - \frac{k T_b}{Q} \ln \left[ \frac{3 * N_{tot} * k * T_0}{2 * \pi * R^3 * P_{atm}} \left( \frac{R_0}{R} \right)^{3(\gamma-1)} \right] \right\}$$

# Fs ablation in a gas : Plume contains NPs

atoms



nanoparticles



**Calculated** plume dynamics for the Ni expansion in  
Ar gas at 300 Pa

up – density of clusters at  $t=10 \mu\text{s}$

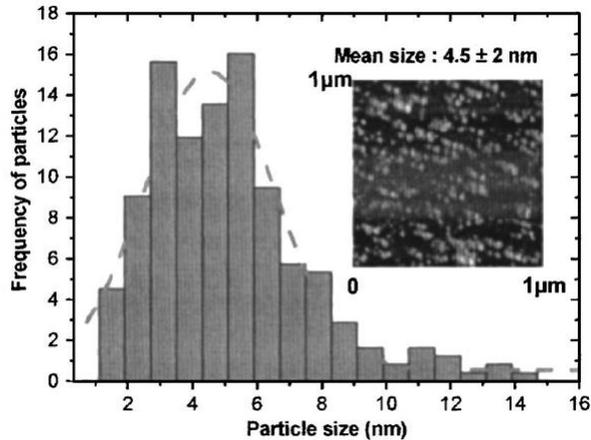
down – density of clusters at  $t=0.55 \mu\text{s}$

**Experiments:**

Amoruso et al. Phys.  
Lett. 93, 191504 (2008)

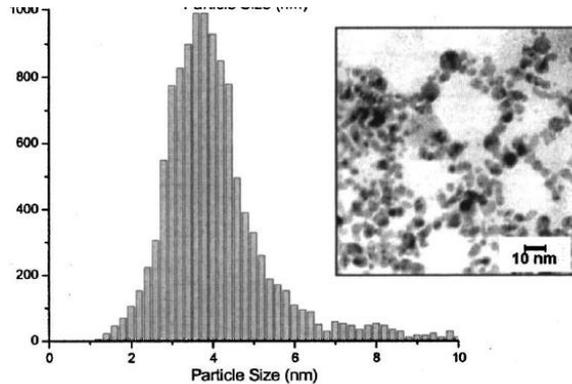
# NP Size Distributions (in a gas or a liquid)

## Laser ablation in air (ns )

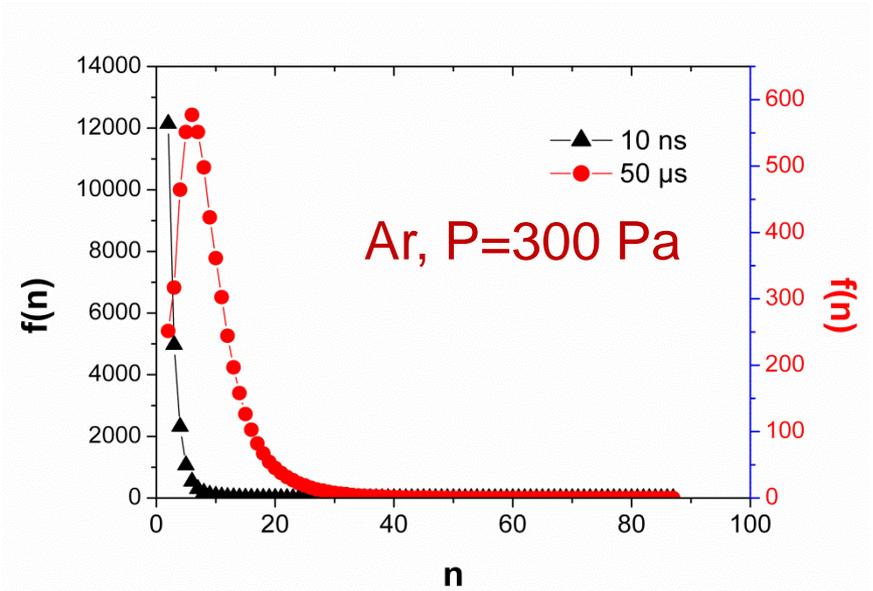


*Pereira et al. AFM steel in air, 25ns*  
 $F=10 \text{ Jcm}^{-2}$

## Fs laser ablation in liquid



*Kabashin et al. Au, TEM  $F=60 \text{ Jcm}^{-2}$*

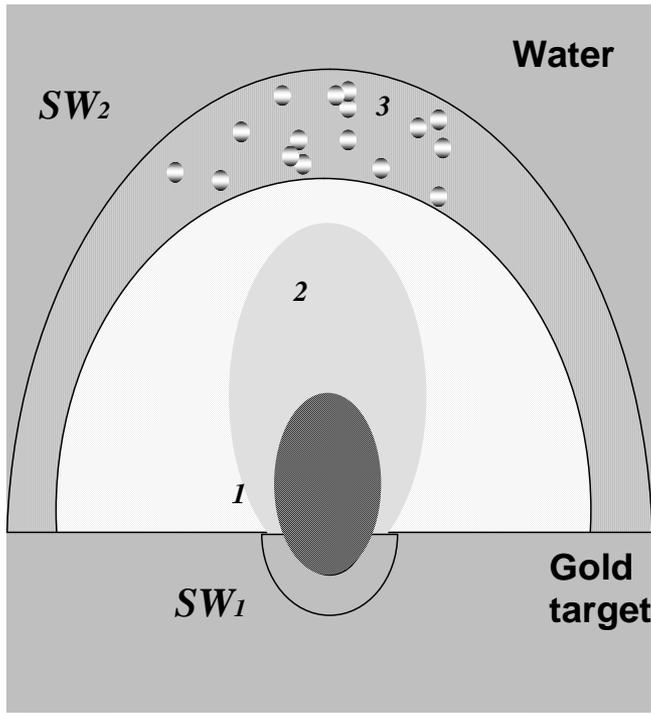


Background gas or liquid  $\Rightarrow$   
LogNormal functions

### ***3. Laser ablation in liquids***

- “Primary” particles: direct ejection, nucleation in bubble or in solution ?
- Secondary” particles: evaporation/growth, coalescence, agglomeration

# NPs formation in liquids



- 1. Early stage** – laser energy absorption, material ejection, plume formation
- 2. Intermediate stage** – plume expansion in the presence of a confining liquid environment, formation of a cavitation bubble (CB)
- 3. Late stage** – CB's collapse, plume mixing with the liquid, NP coalescence/aggregation processes.

## Main questions

- ⇒ **Can NPs be directly ejected from the target?**
- ⇒ **Are NPs formed by condensation in the plume/cavitation bubble or later in liquid solution ?**
- ⇒ **What stage plays the major role ?**

# Ablation in the presence of a liquid

-In a liquid=> different ablation threshold and yield

-shock wave formation, but also

-bubble formation

⇒ Origin of “primary” particles?

⇒ “Secondary” particles will be formed later by coagulation, coalescence/aggregation

⇒ fragmentation

⇒ ...

**Confinement by liquid**

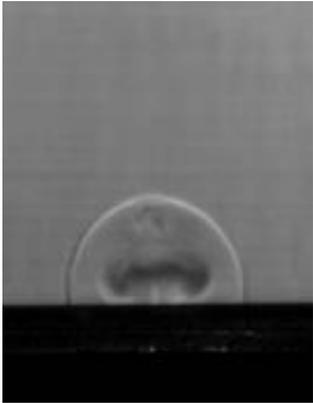
$$R_n(t) = \left\{ R_{na}^3 + \frac{R_{nb}^3 - R_{na}^3}{2\tau} \left[ t - \frac{\tau}{\pi} \sin\left(\frac{\pi}{\tau} t\right) \right] \right\}^{1/3}$$

*Vogel et al., J. Acoust. Soc. Am. 100(1), 1996*

**=>Hydrodynamic modeling**

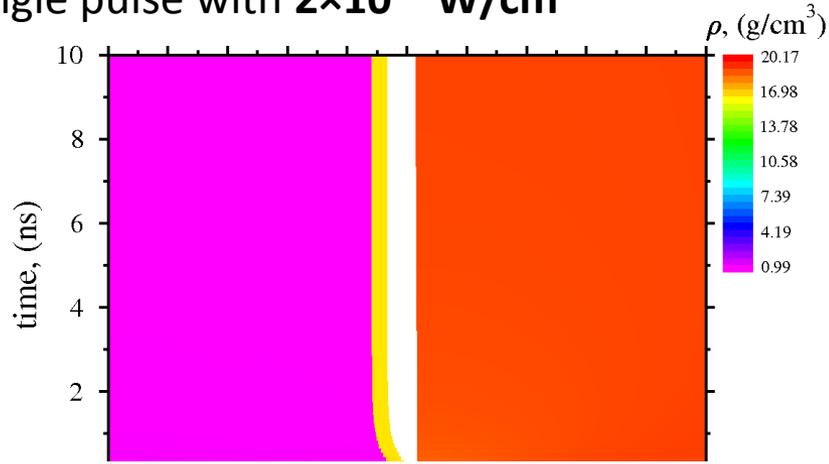
- **Gold target in water**
- *2T semi-empirical equations of state (for metal, gold here)*
- *EOS for water*

$$I(t) = I_0 \exp(-4 \ln 2 t^2 / \tau^2), \quad \tau = 200 \text{ fs}, \quad 800 \text{ nm}, \quad I_0 = 2, 3, 4, 5 \times 10^{13} \text{ W/cm}^2$$

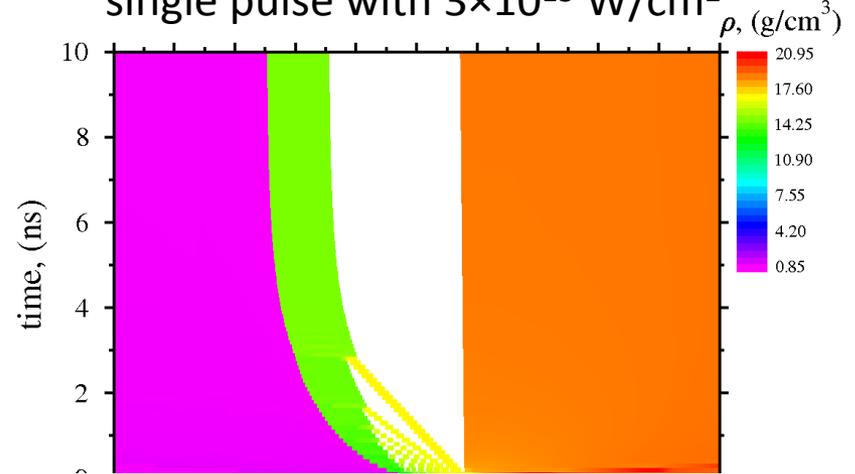


# Single fs pulse, Au in water: simulation close to threshold

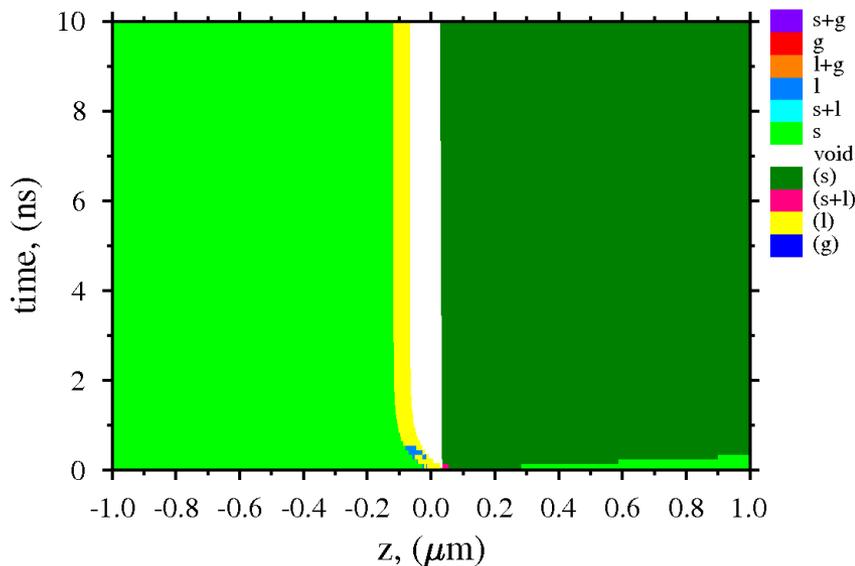
single pulse with  $2 \times 10^{13} \text{ W/cm}^2$



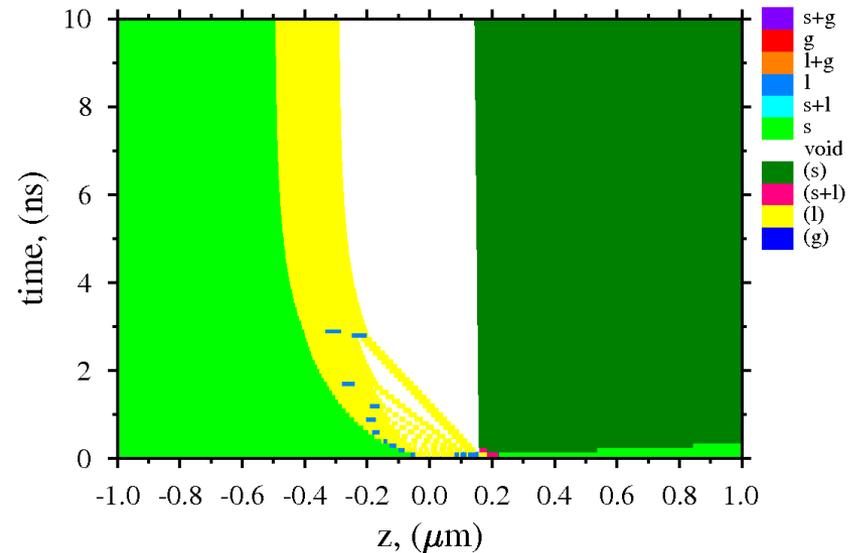
single pulse with  $3 \times 10^{13} \text{ W/cm}^2$



Phase state



Phase state

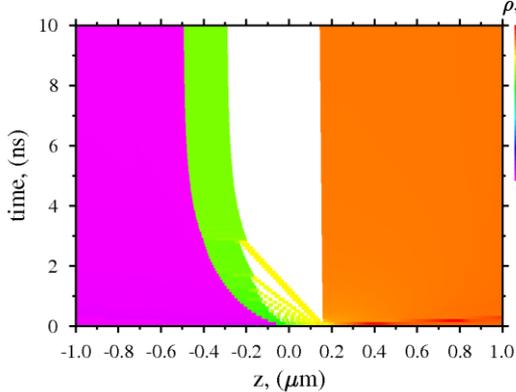


**=> Void formation, higher threshold, fragments cannot escape** 21

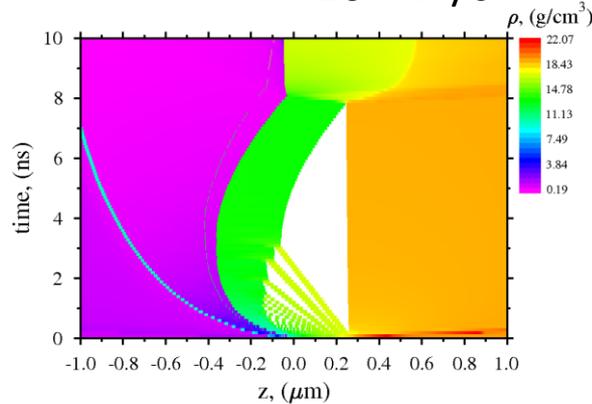
# Increase in laser intensity

Density maps

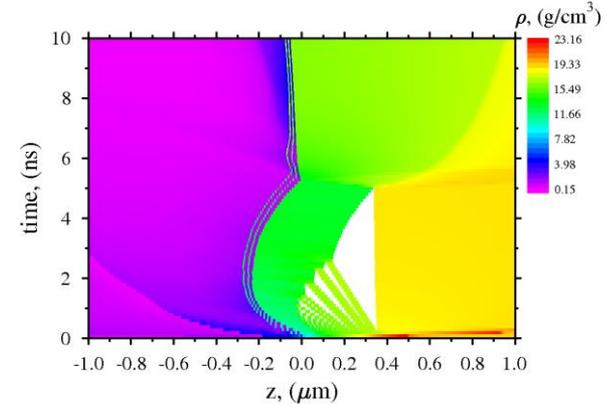
$3 \times 10^{13} \text{ W/cm}^2$



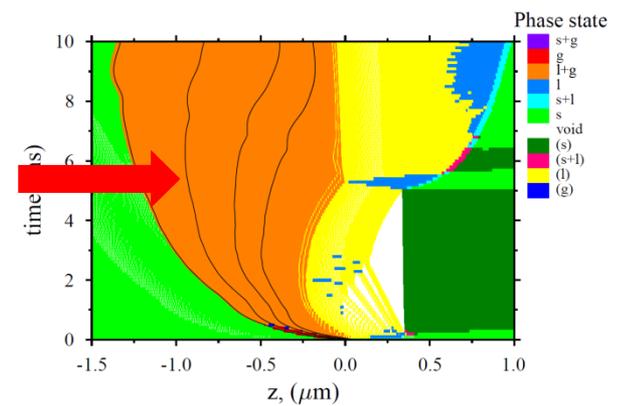
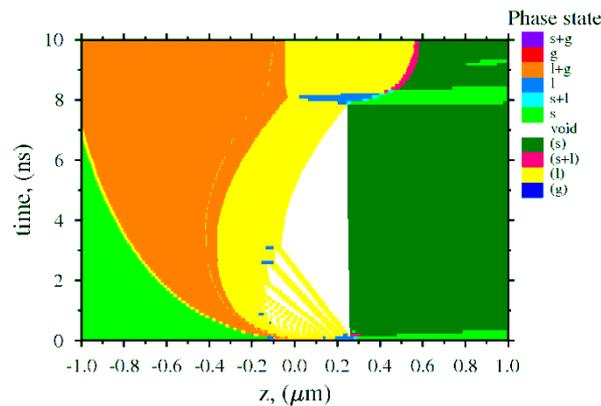
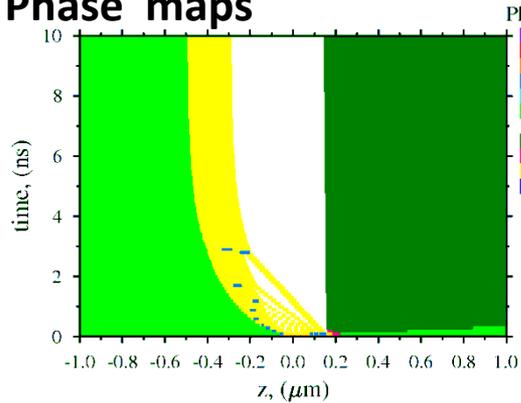
$4 \times 10^{13} \text{ W/cm}^2$



$5 \times 10^{13} \text{ W/cm}^2$



Phase maps

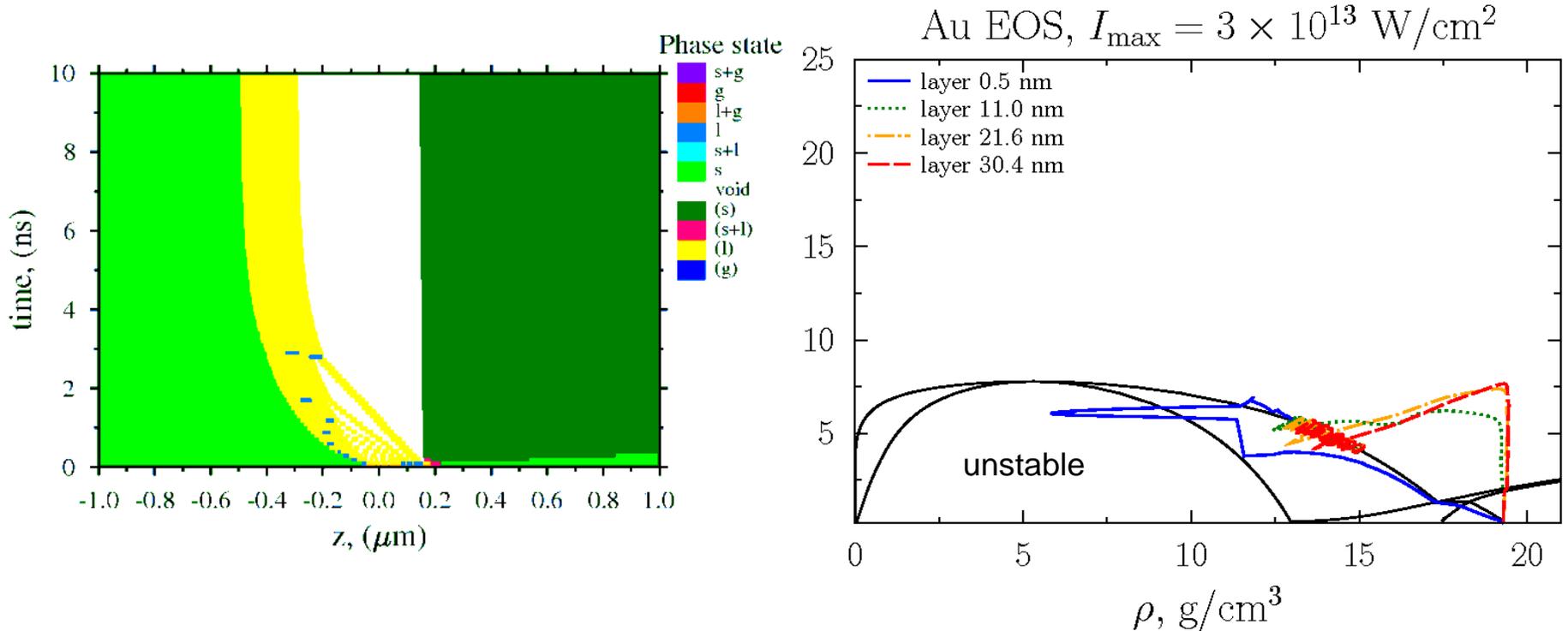


Near ablation threshold, ejected droplets stick together under the pressure

$F \gg$  unstable liquid-gas phase (orange) extends and presses liquid back ( $t=5\text{ns}$ )

**=> NPs are formed in the unstable liquid+gas region (under spinodal) rather than by condensation or fragmentation**

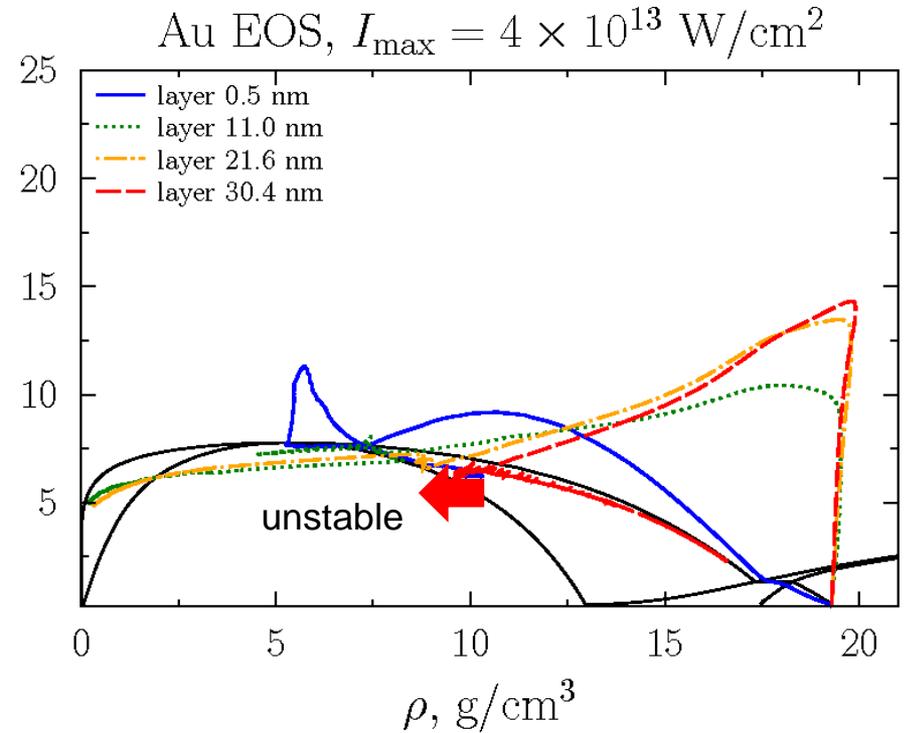
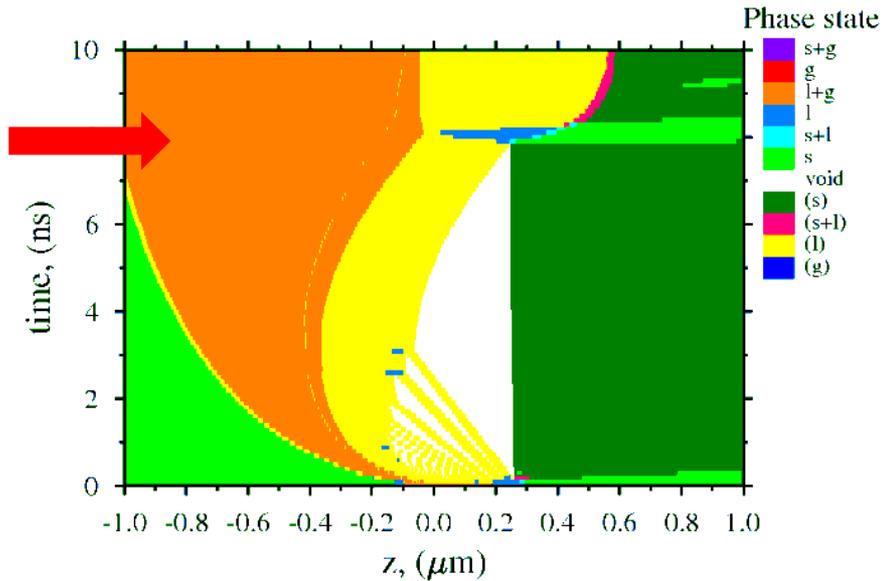
# Near ablation threshold



- Fragmentation does occur, but ambient liquid prevents the expansion
- The fragments stick together under the pressure of the ambient liquid

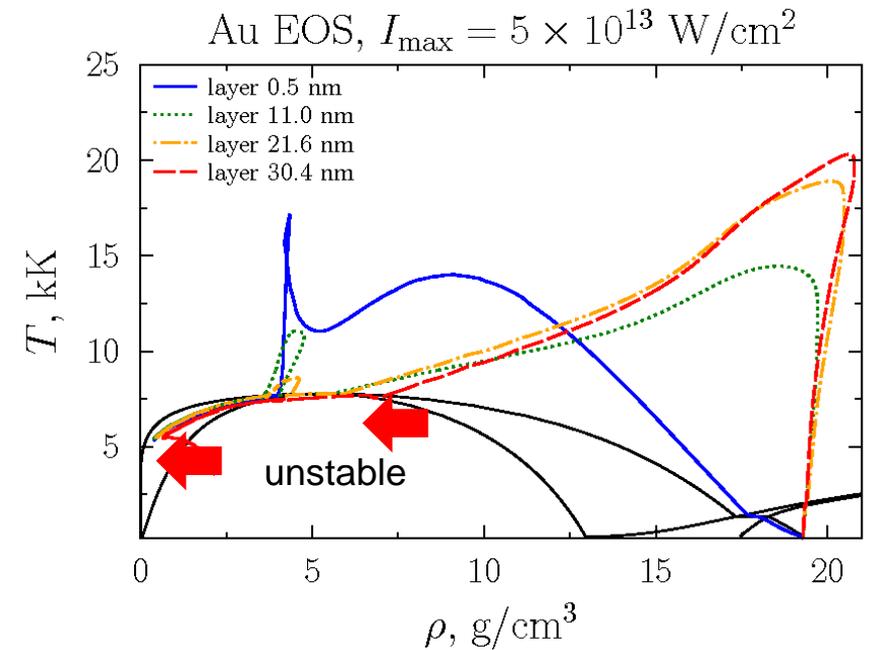
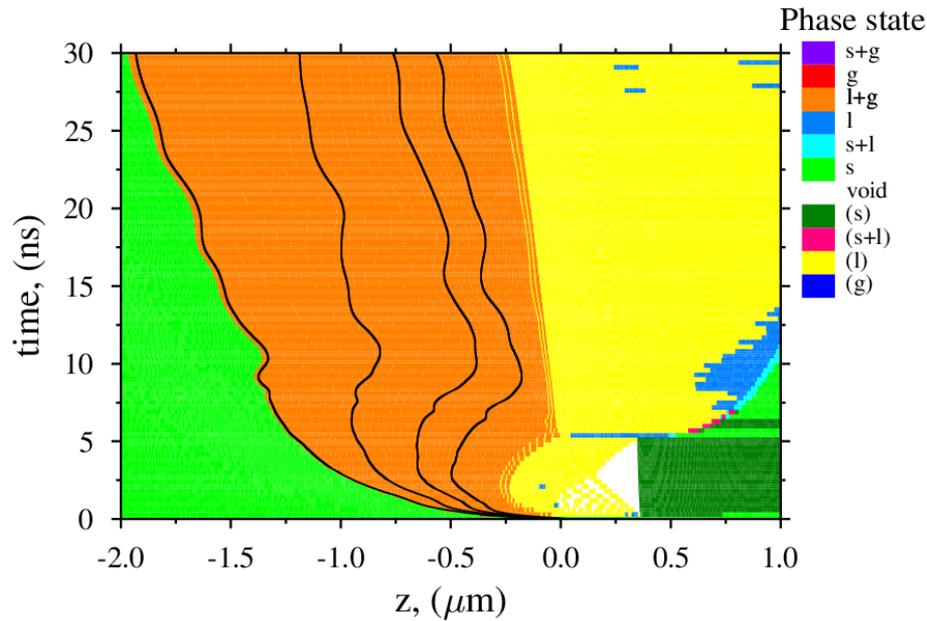
**=> no NPs are expected**

# With the increase in laser fluence



- Upper layers cross the spinodal and enter into unstable liquid-gas region (orange zone)=> **NPs can be formed here**
- Part of trajectories reach supersaturated gas region, where condensation takes place

# Single pulse in liquid: large laser fluence

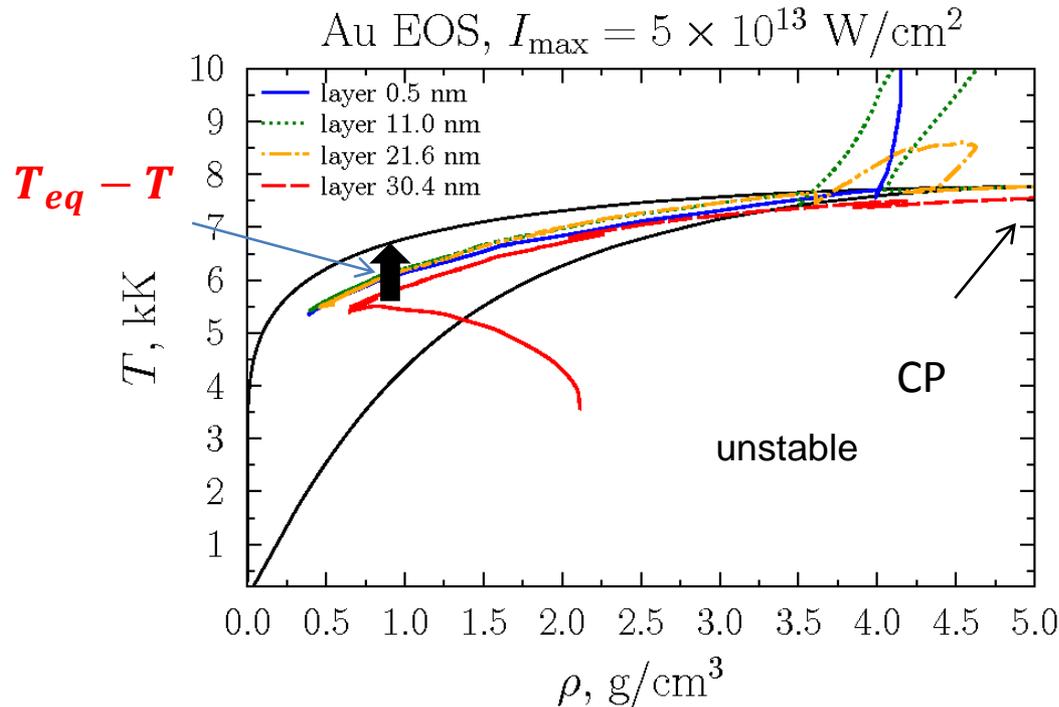


- more trajectories enter to the gas-liquid zone (orange region) => **more important yield of « primary »NPs**
- In addition **condensation** takes place in the region of supersaturated vapor

=> **Two generations of primary NPs:** larger particles from unstable liquid-gas and smaller ones due to nucleation

# Zoom into the region of supersaturation

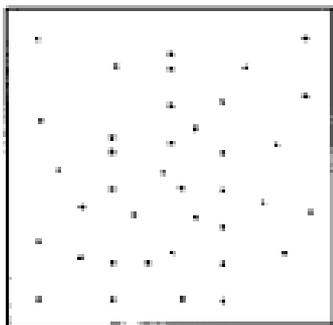
Phase trajectories for 4 layers, single pulse with  $5 \times 10^{13} \text{ W/cm}^2$



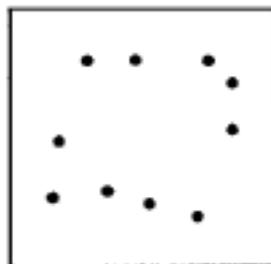
- In the supercooled gas (SCG) region, saturation degree  $\theta = \frac{T_{eq} - T}{T_{eq}} \in \{0.1, 1\}$
- Nucleation, critical radius  $r_c = \frac{2\sigma\omega}{q\theta_{max}} \sim 10^{-8}/\theta_{max} \sim 10 - 100 \text{ nm}$

# Later stage: two processes

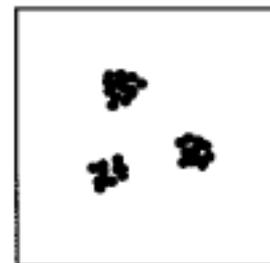
Supersaturated solution



**1-nucleation**  
In liquid solution



**2-coalescence/aggregation**

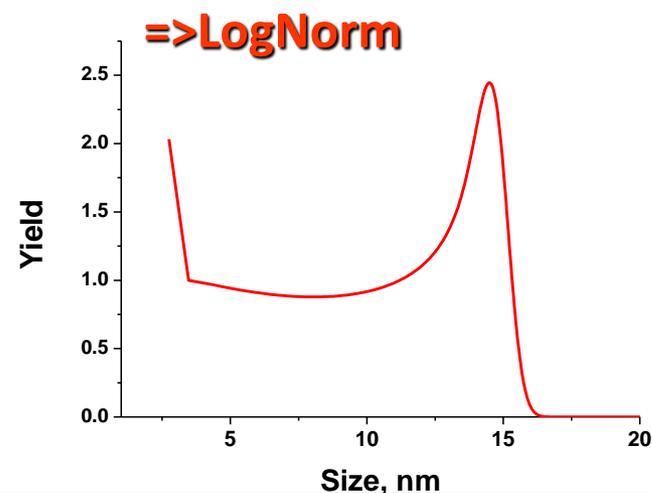


$$\rho(t) = K_c c^2 \exp\left[\frac{-\Delta G(n_c, c)}{kT}\right]$$

$$\frac{dN_1}{dt} = \rho(t) - \sum_{j=2}^{\infty} j \frac{dN_j}{dt} \quad \frac{dN_2}{dt} = fK_1 N_1^2 - K_2 N_1 N_2$$

$$\frac{dN_s}{dt} = K_{s-1} N_1 N_{s-1} - K_s N_1 N_s \quad (s \geq 3)$$

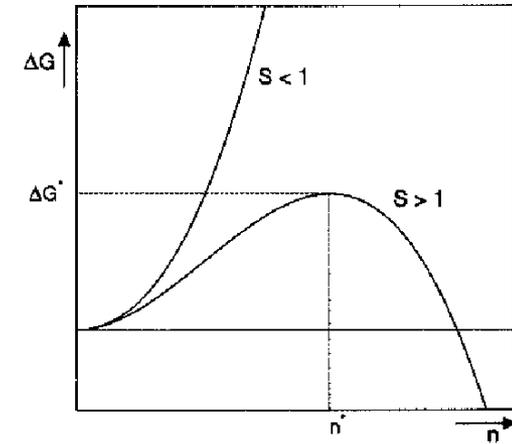
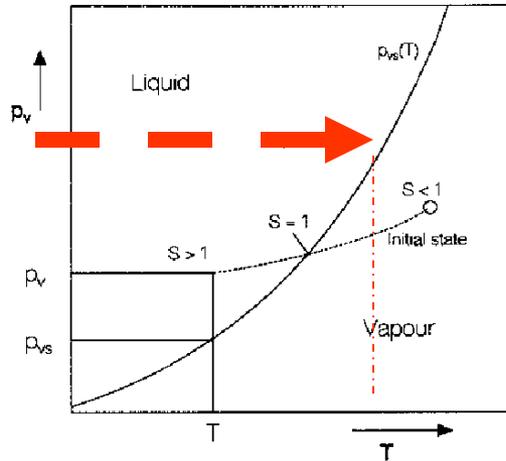
$$K_n = 4\pi(a + an^{1/3})(D_a + D_a n^{-1/3}) \approx 4\pi a n^{1/3} D_a$$



# Diffusion-driven nucleation

$$\Delta G(n, c) = -nkT \ln(c/c_0) + 4\pi a^2 n^{2/3} \sigma$$

## Saturation curve



## Supersaturation ratio

$$S = c / c_0$$

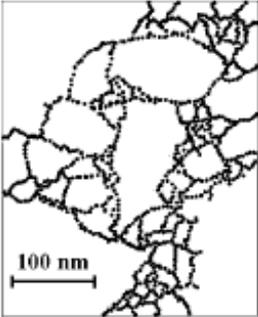
The peak of the nucleation barrier corresponds to the critical cluster size

$$n_c = \left[ \frac{8\pi a^2 \sigma}{3kT \ln S} \right]^3$$

$$J_{nuc} \sim \exp \left[ \frac{-\Delta G(n_c, c)}{kT} \right]$$

At given T pressure should be high => need of high loading

# Secondary particles => Aggregation/Coalescence



## Smoluchowski master equation

$$\frac{\partial n}{\partial t} = \frac{1}{2} \int_0^v \beta(\tilde{v}, v - \tilde{v}) n(\tilde{v}) n(v - \tilde{v}) d\tilde{v} - \int_0^\infty \beta(v, \tilde{v}) n(\tilde{v}) n(v) d\tilde{v}$$

$$N_p = \frac{v}{v_p} = A \left( \frac{r}{r_p} \right)^{D_f}$$

Petrus J. Dekkers<sup>\*,1</sup> and Sheldon K. Friedlander

### Kernel depends on the regime

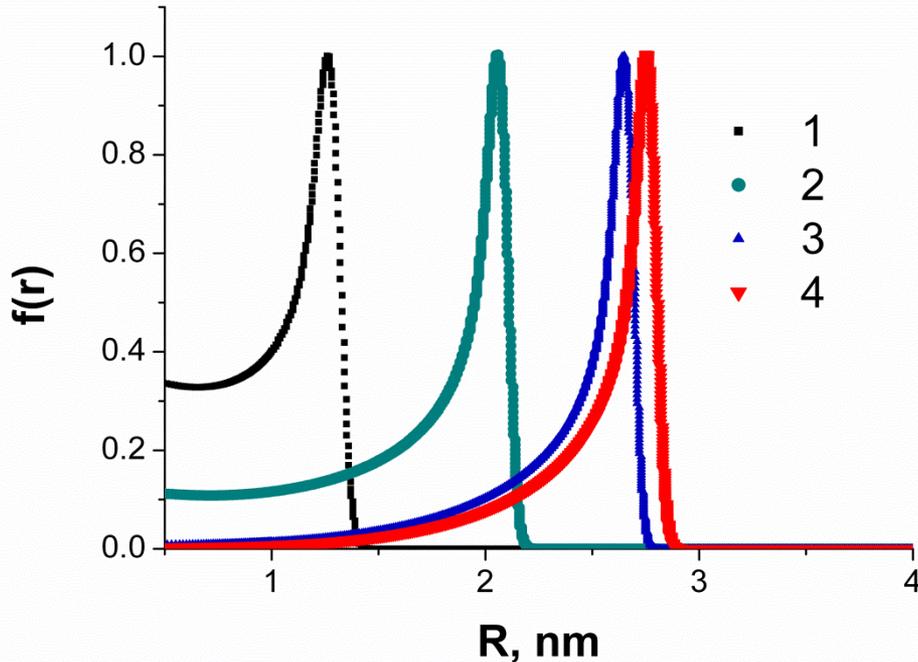
1. Free molecular regime ( $Kn \gg 1$ )
  2. Continuum regime ( $Kn \ll 1$ )
  3. Near-Continuum transition regime ( $0.01 < Kn < 1$ )
- **The distribution is the narrowest in the near-continuum transition regime ( $0.01 < Kn < 1$ ) when there are collisions between the primary particles, but not too much**

# Collisional growth of secondary particles

$$\frac{dN_1}{dt} = \rho(t) - \sum_{j=2}^{\infty} j \frac{dN_j}{dt}$$

$$\frac{dN_2}{dt} = fK_1N_1^2 - K_2N_1N_2$$

$$\frac{dN_s}{dt} = K_{s-1}N_1N_{s-1} - K_sN_1N_s \quad (s \geq 3)$$



Calculated size distribution obtained for 1-10, 2-100, 3-1000 and 4-2000 pulses. Here, laser frequency is 1 kHz, gold solution in water is considered with  $a=1.9 \cdot 10^{-10}$  m.

# How to reduce mean size and/or size dispersion ?

-Reduce the number of collisions to avoid aggregation

Ablation yield/loading reduction=>near free-molecular regime

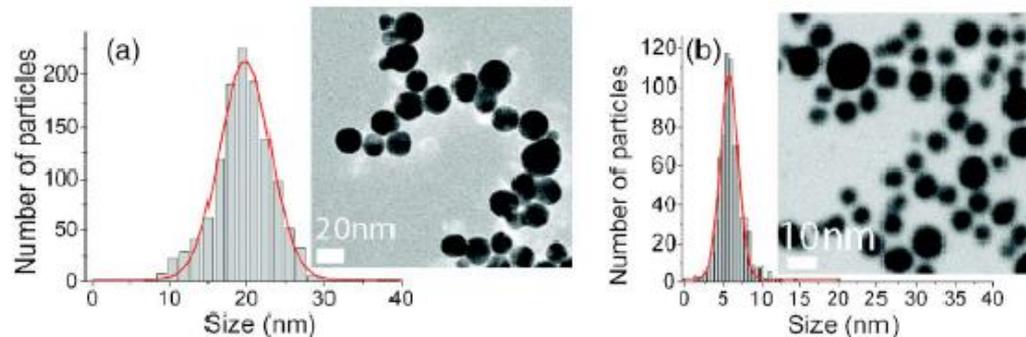
-Reduce the probability of aggregation

Chemical composition of the solution, surfactant molecules

-Reduce background temperature

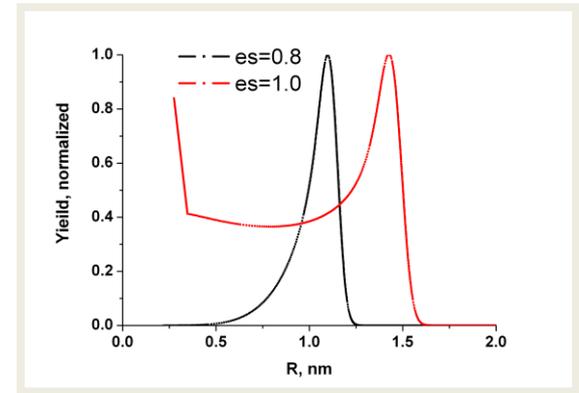
-Irradiate solution by light

(continuum generation followed by particle fragmentation in a solution)



# How to reduce NP's size ?

⇒ **Surfactant molecules**  
Change in stiking parameter

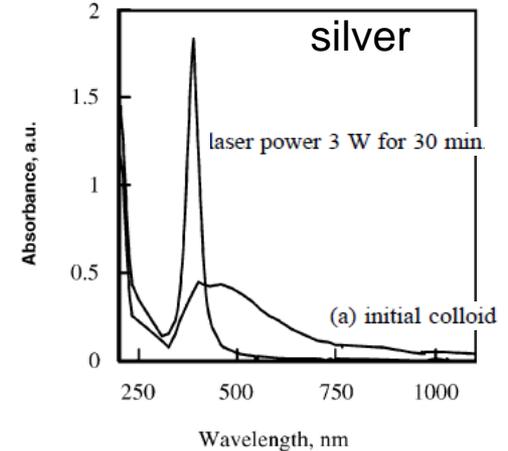
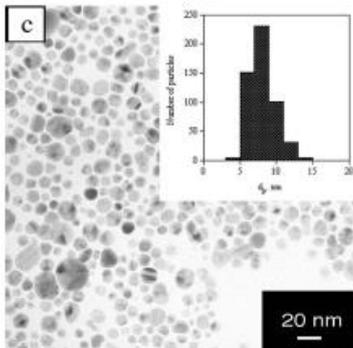
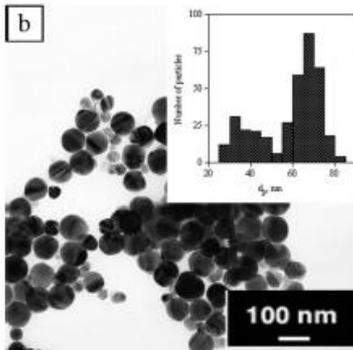
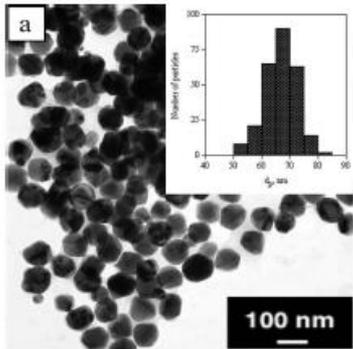


⇒ **Additional laser irradiation**

depends on laser  
wavelength, fluence and  
size of particles

⇒ **mechanisms ?**

*Träger et al.*  
*Sylvestre et al. 2004*  
*Pustovalov et al. 2005*  
*Pyatenko et al. 2005*

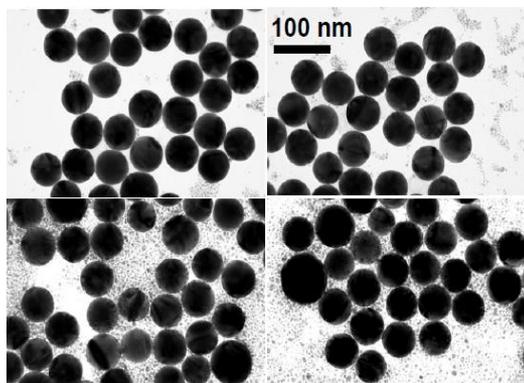


## ***4. Laser-induced particle size reduction and fragmentation***

# MOTIVATION

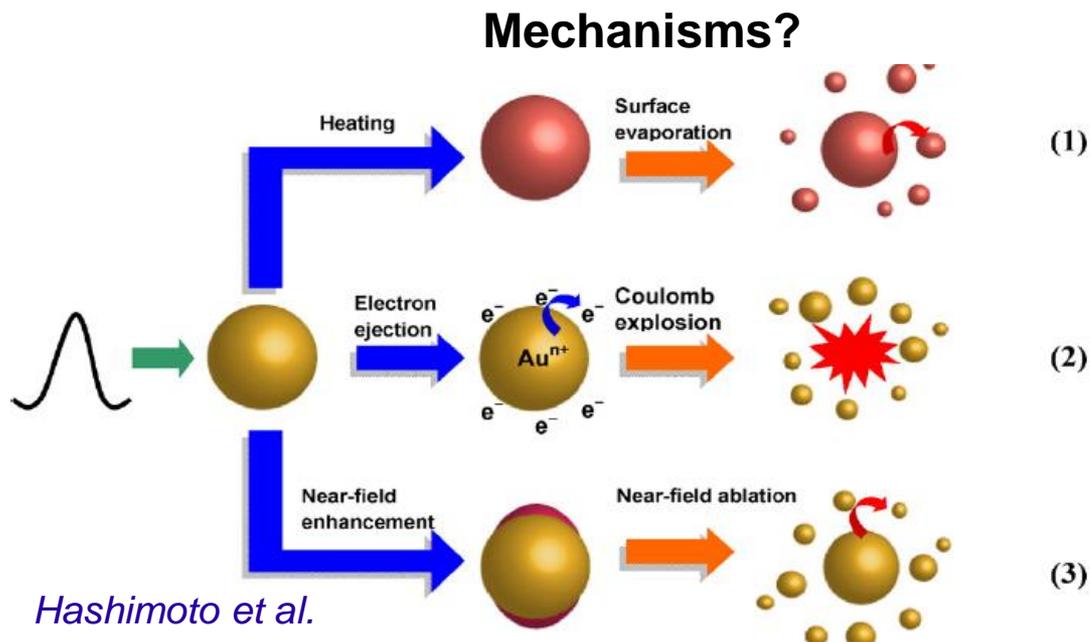
**Main goal: *understanding ultra-short laser interactions with metallic nanoparticle (NP) for better control over laser-assisted NP formation in liquids***

Werner et al.  
*J. Phys. Chem. C*,  
2011, 115 (12), 5063



*J. Phys. Chem. C*,  
2011, 115 (12),  
8503

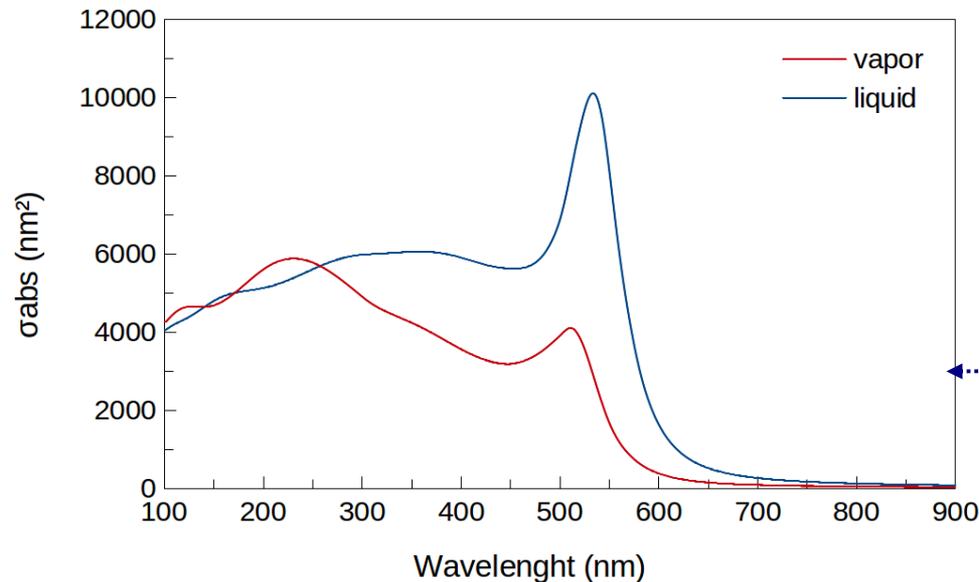
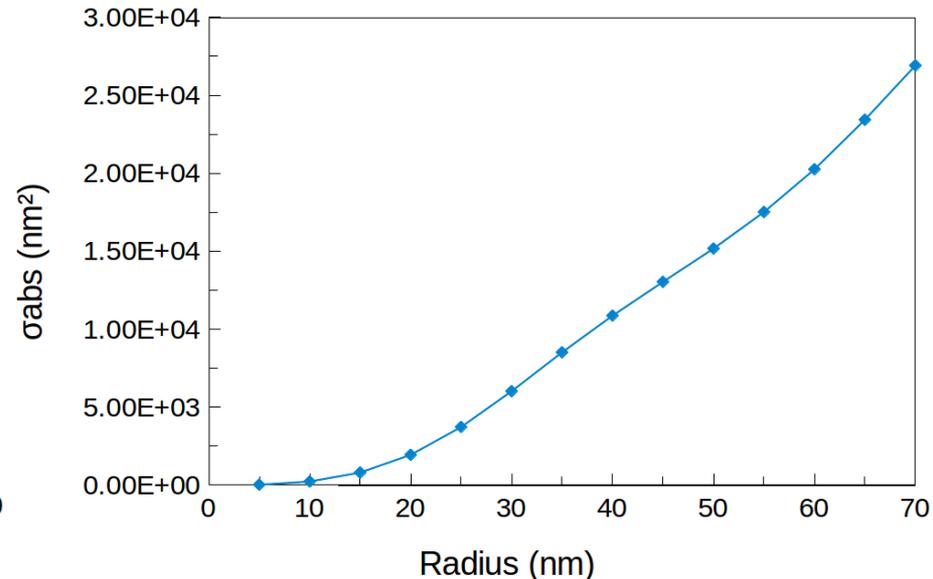
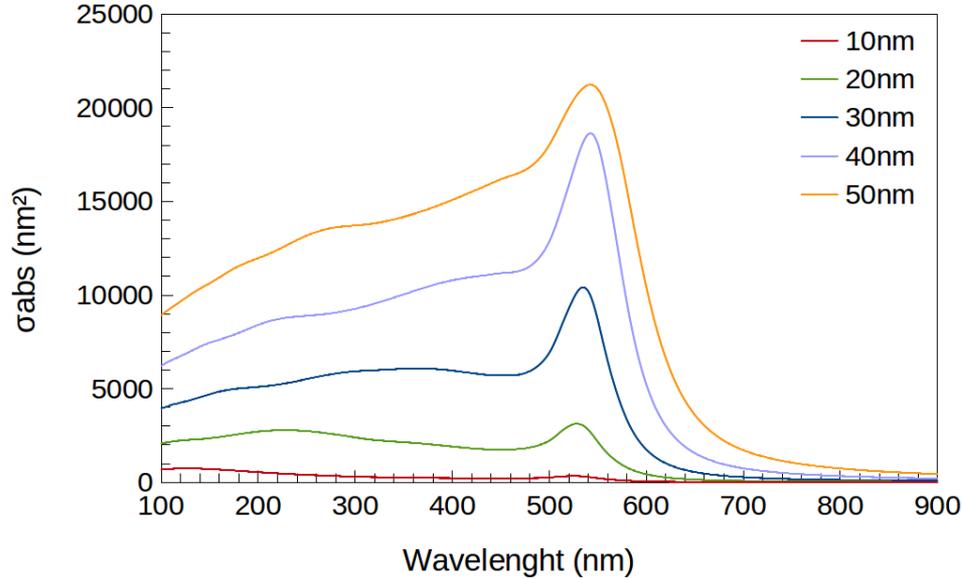
For different applications, such as  
*-photodynamic therapy, cancer treatment*  
*-photonic and plasmonic devices*  
*-solar cells, novel energy sources*  
*-sensors*



*Journal of Photochemistry and Photobiology C: Photochemistry Reviews* 13 (2012) 28–54

# ABSORPTION

## Au NPs with radius in the range 10-100 nm

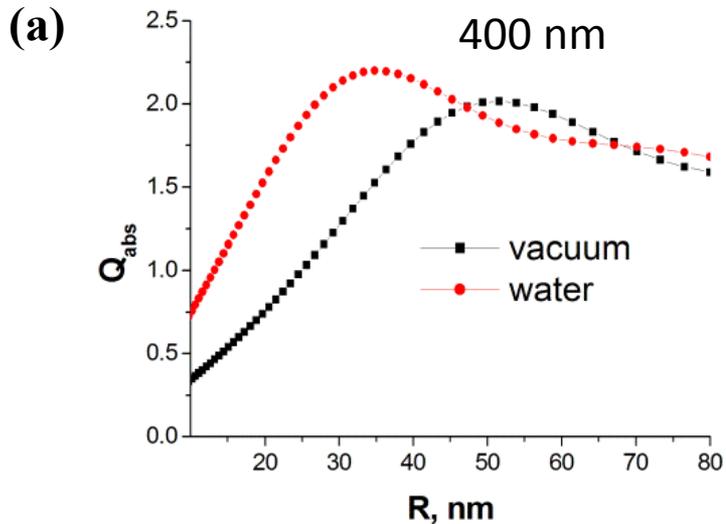


Absorption cross section in liquid water raises with the particle radius (not linear) at 400 nm.

Absorption depends of the medium. At 400 nm, a change from liquid water to vapor results in a considerable decay in absorption.

# ABSORPTION COEFFICIENT vs RADIUS

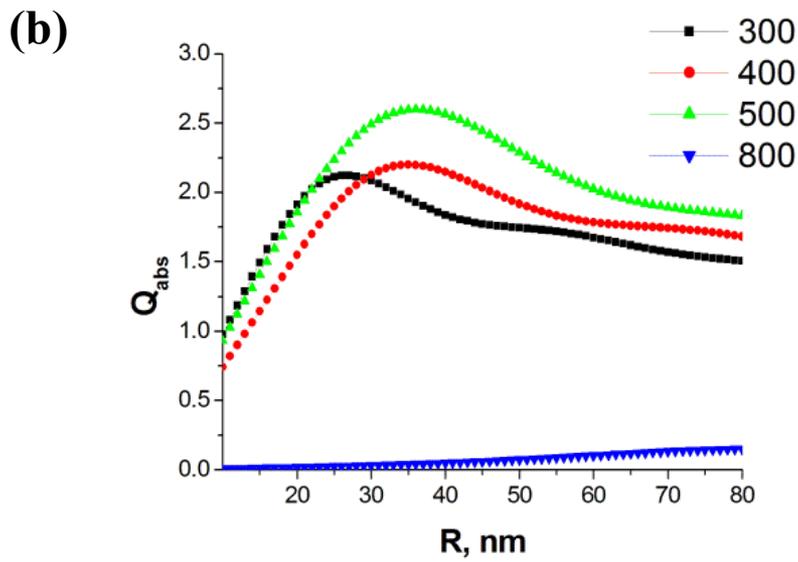
## Gold NPs water



**Generalized Mie theory** =>

1- calculated absorption coefficients vs AuNP's radius,  $R$  For water (red) and vacuum (black) at 400 nm

**=> peak in absorption at  $R \sim 30$  nm for AuNP in water**



2- calculated absorption coefficients vs NP's radius,  $R$  for AuNPs in water and different wavelength (300-800nm)

**=> peak position is shifted to larger  $R$  if wavelength rises**

# HEATING by FS LASER: TWO TEMPERATURE MODEL (TTM)

Laser => heating of a metallic particle  
energy absorbed by the conduction electrons

Energy is transferred :  
- first to the lattice (electron-phonon collision)  
- then to the surrounding medium (thermal diffusion)

solve

$$\left\{ \begin{array}{l} C_e(T_e).dT_e/dt = -g(T_e).(T_e-T_l) + S(t) \\ C_l(T_l).dT_l/dt = g(T_e).(T_e-T_l) - F \\ C_m(T_m).dT_m/dt = F \end{array} \right.$$

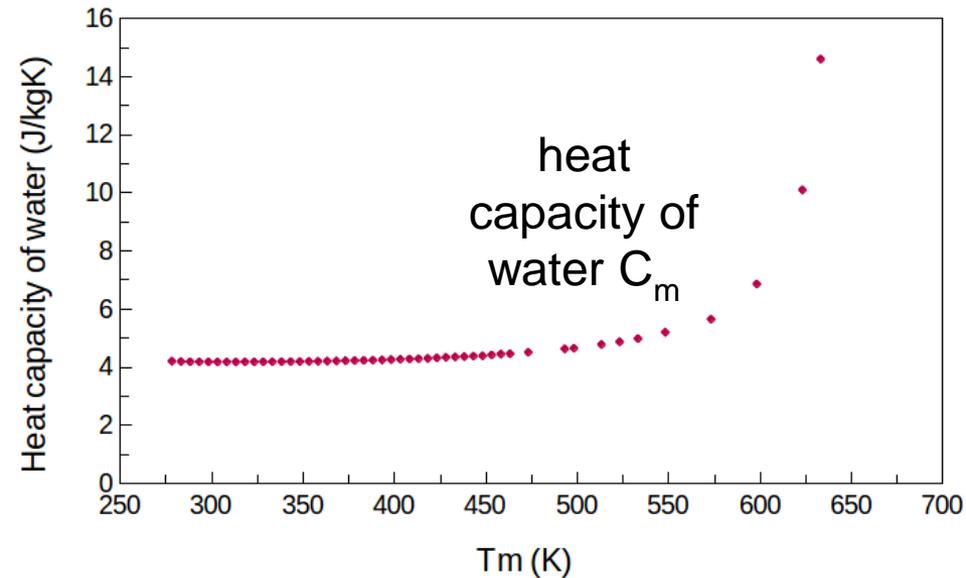
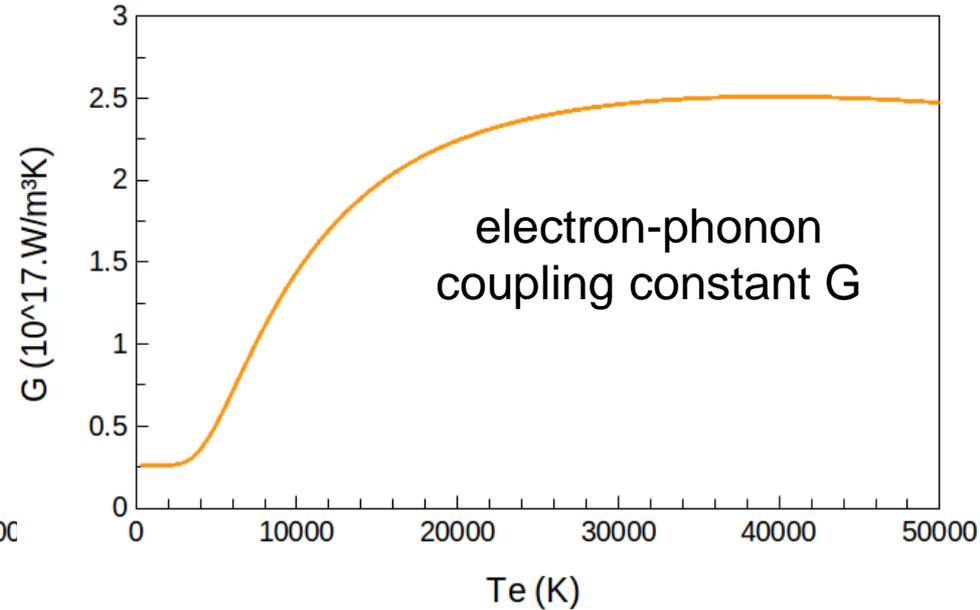
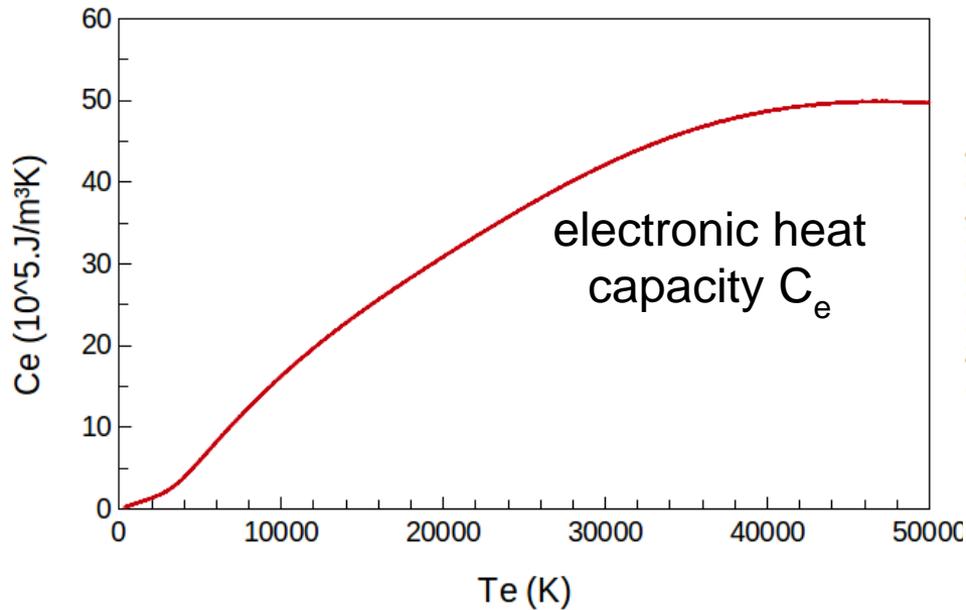
*evolution of  $T_e$  coupling to the lattice*  
*dynamics of  $T_l$  including heat transfer*  
*neglect heat conduction during calculation time ~ several ps*

with

$$\left\{ \begin{array}{l} S(t) = \sigma_{mie}(n_m, R_{NP}).P(t) / V(T_l) \\ P(t) = (2\sqrt{\ln 2} \cdot \text{flu} / \tau \sqrt{\pi}) * \exp(-4(\ln 2)t^2 / \tau^2) \\ F = 3h.(T_l - T_m) / R_{NP} \end{array} \right.$$

*S(t) energy-exchange with the laser pulse*  
*P(t) gaussian pulse*  
*F empirically found energy loss term*

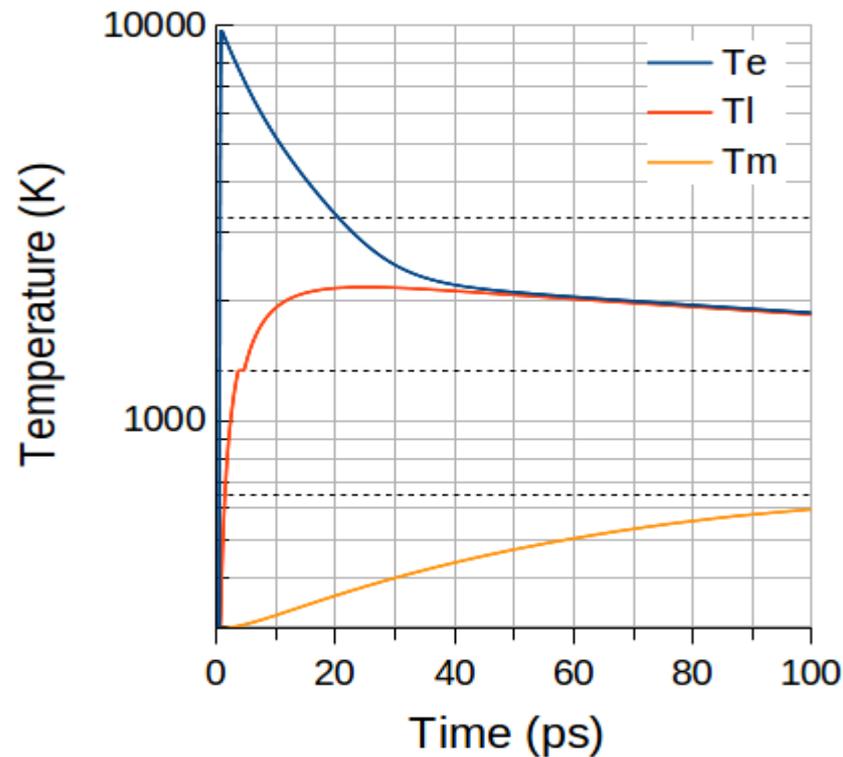
# PARAMETERS FOR TTM



=> based on electron density of states, taking into account the effect of thermal excitation of the electrons (including  $d$  electrons) located below  $E_F$

Lin et al.  
*Phys.Rev. B* 77, 075133 (2008)

# ELECTRONIC & LATTICE TEMPERATURES



**30 nm** radius gold particle in water  
absorbing laser pulse of **150 fs** at **400 nm**  
with laser fluence of **12.3 mJ/cm<sup>2</sup>**  
Gaussian time profile

Agreement with

*Werner et al. J. Phys. Chem. C, 2011, 115 (17), 8503*

=>  $T_e$  rises instantaneously, then energy deposited into the electronic system is transferred to the lattice =>  $T_l$  rises

=> Particle cools through heat exchange with the surrounding water

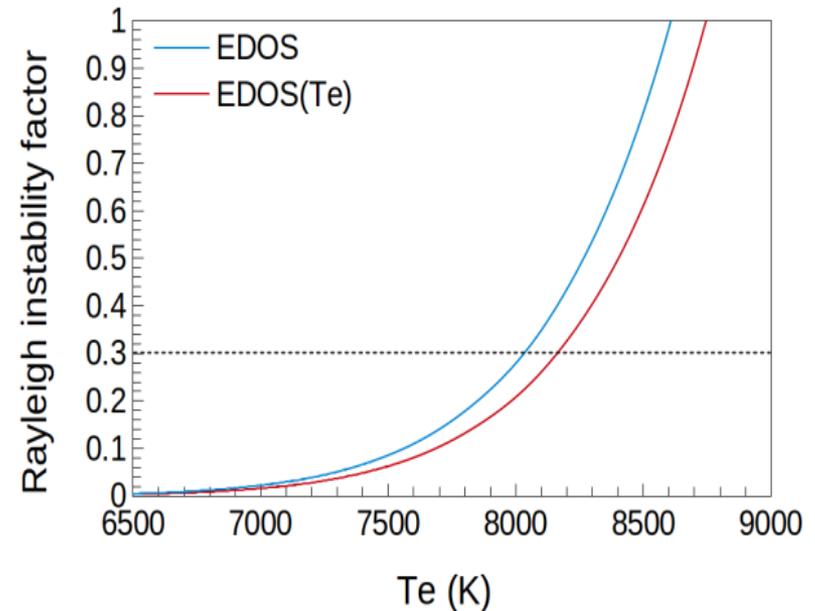
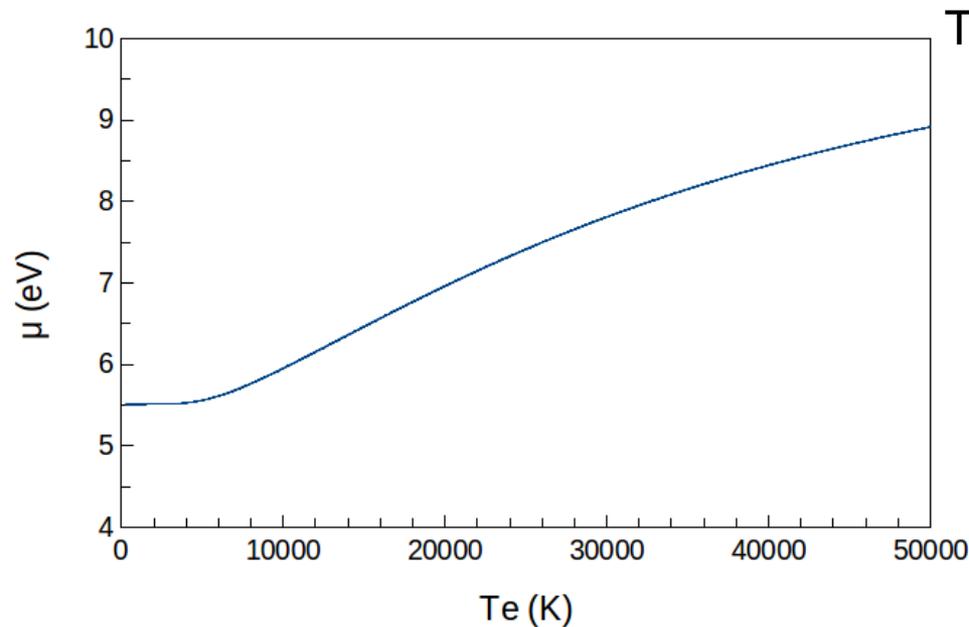
# ELECTRON EMISSION

$N_{\text{therm}}$  is the total **number of emissible electrons** :  $N_{\text{therm}} = \alpha \cdot n_{\varepsilon}$

where  $n_{\varepsilon}$  is a number of electrons with kinetic energy that exceeds the work function per atom at an electron temperature of  $T_e$  :

$$n_{\varepsilon} = \int_{\varepsilon}^{\infty} \text{EDOS}(E) \cdot f(E, \mu(T_e), T_e) \cdot dE$$

electronic density of states for a broad range of  $T_e$  chemical potential value depends on



*Bévillon et al. Phys. Rev. B 89, 115117 (2014)*

**The criterion is based on Rayleigh instability factor (“Liquid drop”):**

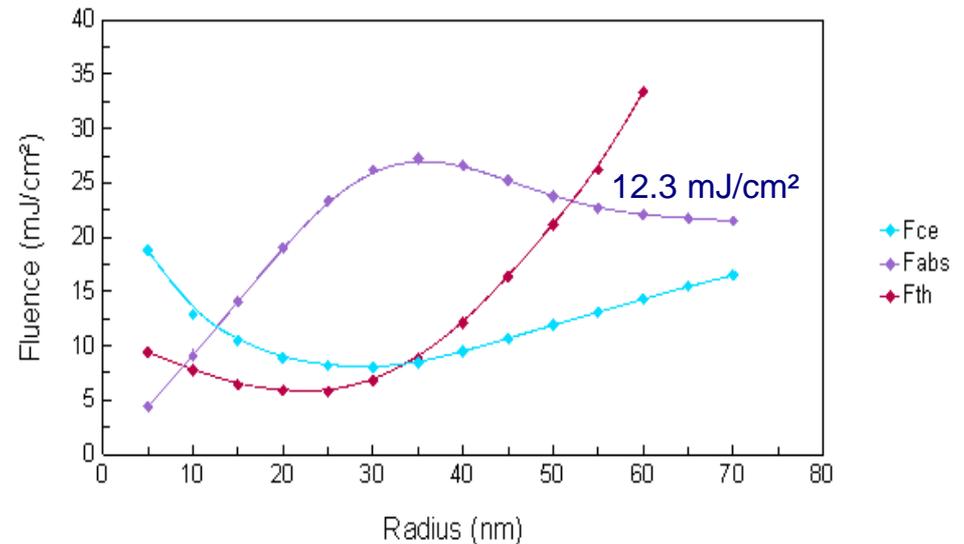
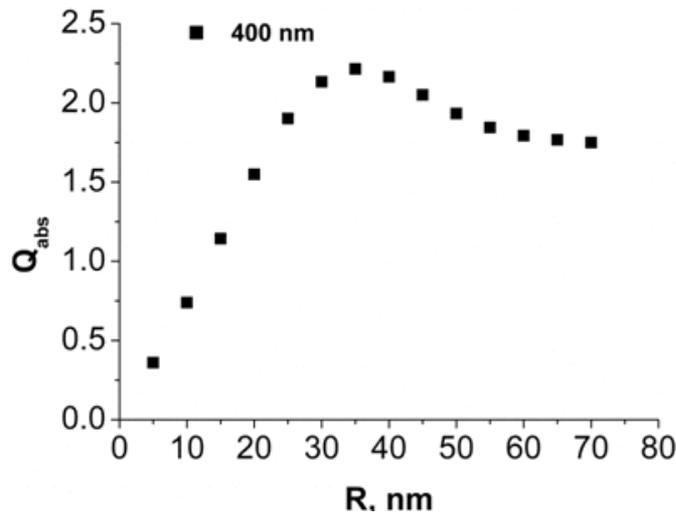
Electrostatic explosion can occur **if  $X > 0.3$** , where  $X = (N_{\text{therm}}^2 / N_e) / (16 \cdot \pi \cdot r_{\text{ws}}^3 \cdot \sigma / e^2)$

$$\alpha = 4V / a_{\text{fcc}}^3 \text{ total number of atoms in the particle}$$

# RESULTS for fs laser interaction with AuNPs

## Absorption coefficient vs radius

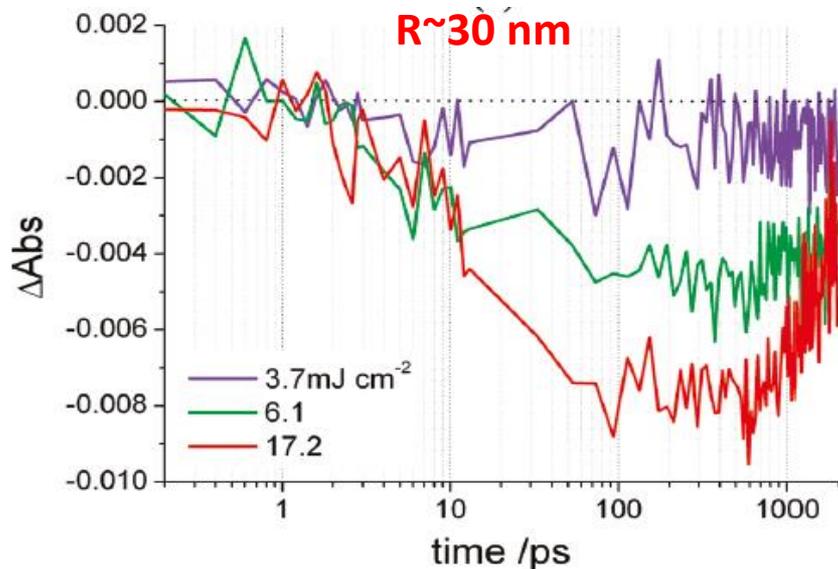
Absorbed laser fluence,  
fluence required for melting,  
fluence required for Coulomb explosion  
for  
150 fs at 400 nm



Fluences required for melting are smaller than the ones  
for CE if radius  $< \sim 30$  nm

# RESULTS FOR fs LASER and AuNPs

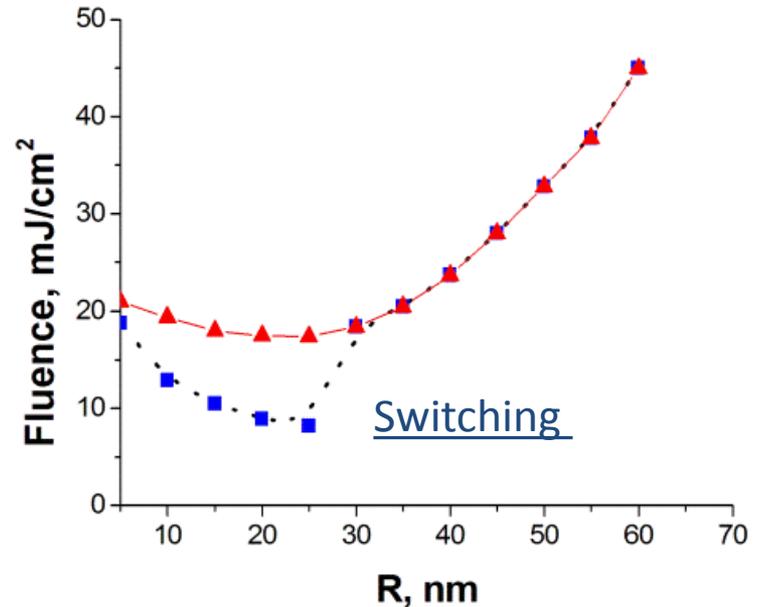
## Experiments



Werner et al. *J. Phys. Chem. C*, 2011, 115, 8503

Threshold fluence

## Modelling



**If  $F > \sim 17.2$  mJ/cm<sup>2</sup>, thermal evaporation of  $R \sim 30$  nm AuNP occurs**

**For smaller  $F \Rightarrow$  no size reduction of 30 nm particles**

$\Rightarrow$  To decompose such particles, one needs an increased fluence

$\Rightarrow$  confirms the switching to the red curve

## 5.1 Summary

- *Primary NPs* are formed in or even before the cavitation bubble **mostly by spinodal decomposition and by condensation**
- **Later stage => formation of “*secondary*” particles** due to nucleation and aggregation
- **Size distribution is log-normal** *due to collisional growth*
- **Mean size of nanoparticles can be also reduced by**
  - Surfactant molecules
  - Additional laser irradiation



## Discussions with:

*S. Mottin,  
A. Tishchenko*

## Support from:

*PALSE ERTIGO project  
PICS CNRS/RAS project  
FP7 EU project BUONAPART-E  
CINES of France*

**THANK YOU FOR YOUR ATTENTION!**

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